

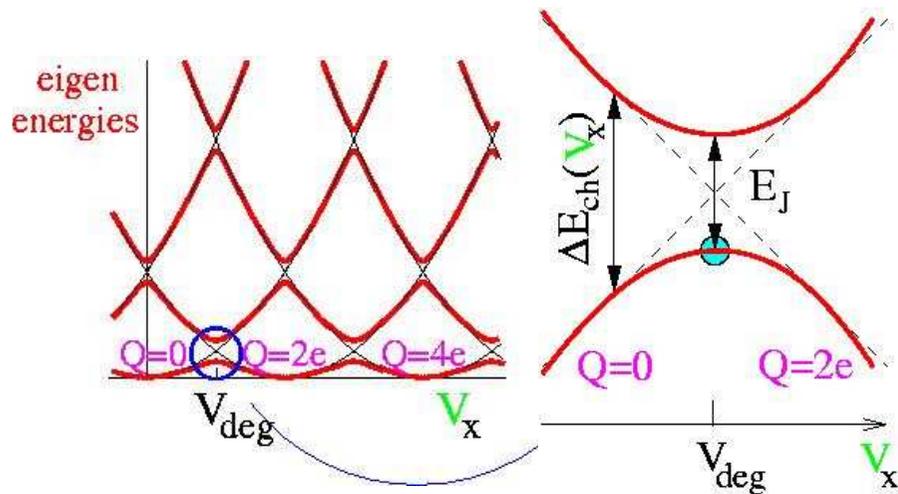
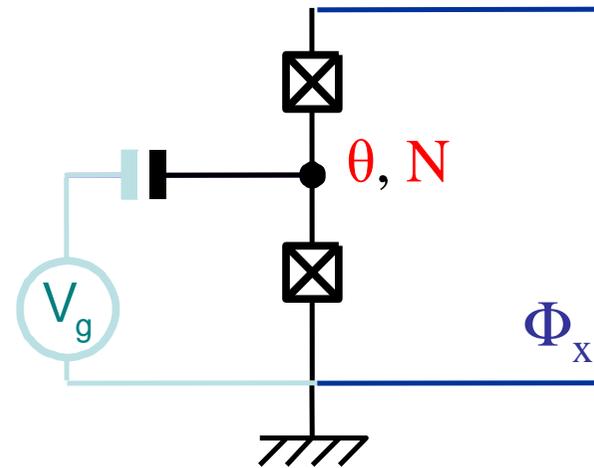
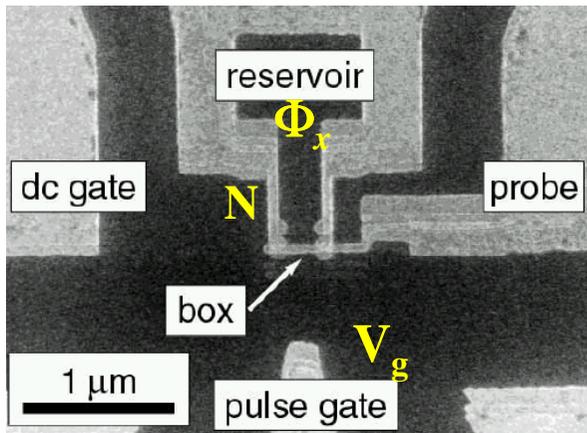
Period-doubling bifurcation readout for a Josephson qubit

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- quantum readout
- parametric driving and bifurcation
- stationary states, switching curve
- discussion

Josephson charge qubits



$$H = E_C \left(N - \frac{C_g V_g}{2e} \right)^2 - E_J \cos\left(\pi \frac{\Phi_x}{\Phi_0}\right) \cos \theta$$

tunable E_J

↓ 2 states only, e.g. for $E_C \gg E_J$

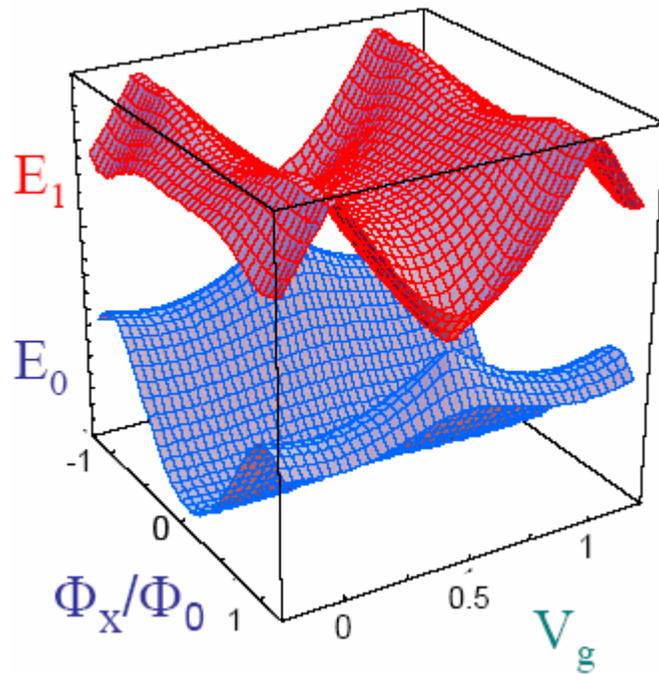
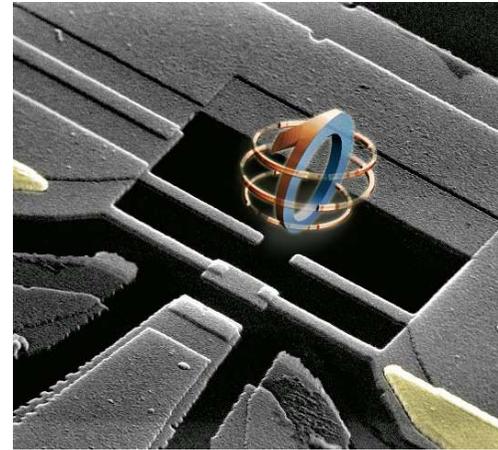
$$H = -\frac{1}{2} E_{ch}(V_g) \sigma_z - \frac{1}{2} E_J(\Phi_x) \sigma_x$$

Charge-phase qubit

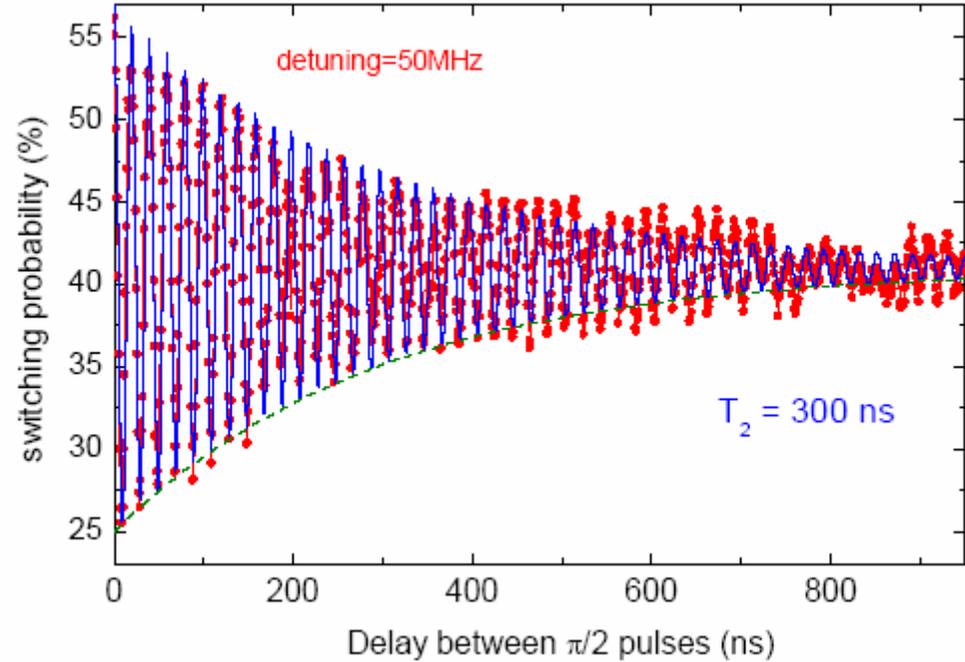
Vion et al. (Saclay)

$$H = -\frac{1}{2} E_{ch}(V_g) \sigma_z - \frac{1}{2} E_J(\Phi_x) \sigma_x$$

Quantrium

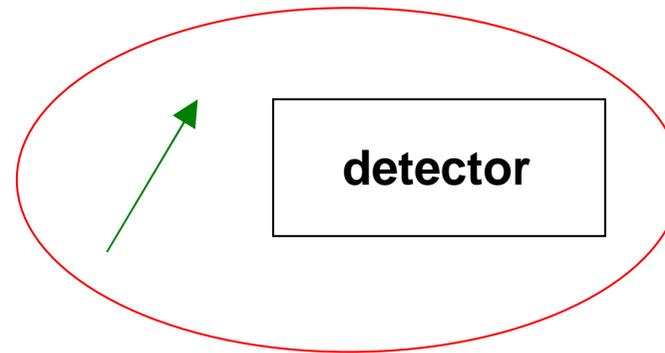


operation at a saddle point



Quantum measurement

$$a|0\rangle + b|1\rangle$$



quality of detection:

- reliable ? single-shot?
- QND? back-action
- fast?

result	0	1
state	$ 0\rangle$	$ 1\rangle$
<hr/>		
probability	$ a ^2$	$ b ^2$

Measurement as entanglement

unitary evolution of qubit + detector

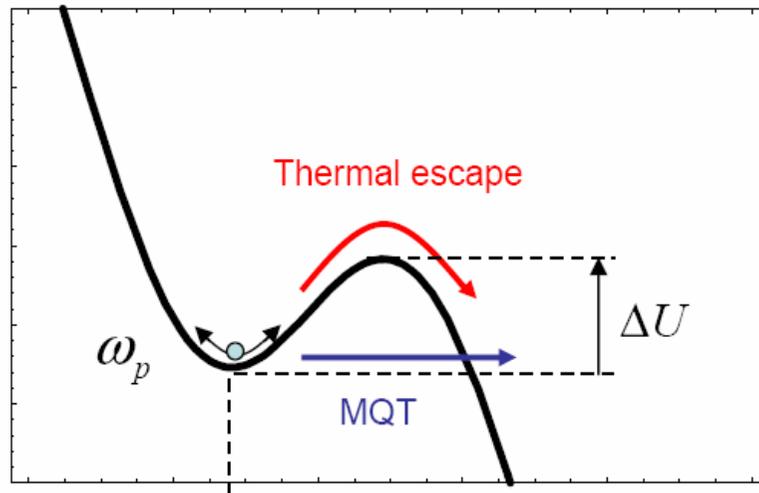
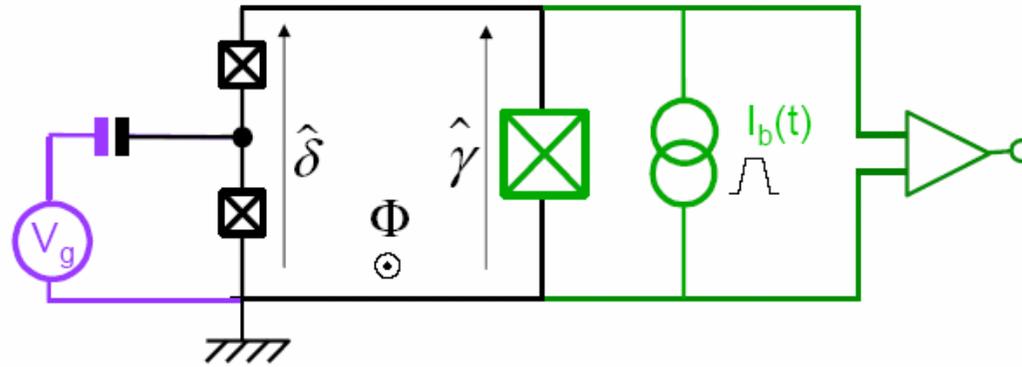
$$(a|0\rangle + b|1\rangle) \otimes |M\rangle$$

\Downarrow

$$a|0\rangle \otimes |M_0\rangle + b|1\rangle \otimes |M_1\rangle$$

$|M_i\rangle$ --- macroscopically distinct states

Quantum readout for Josephson qubits



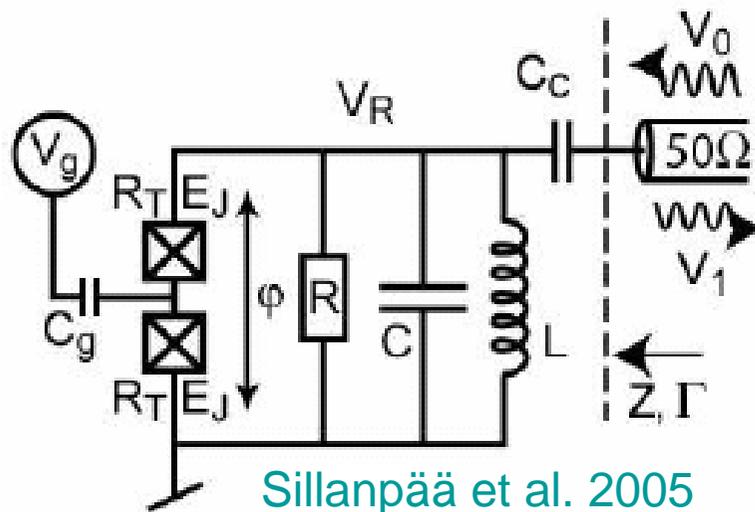
$$U = -I_c \cos\varphi - I \varphi$$

switching readout – monitor the critical current

- no signal at optimal point
- strong back-action by voltage pulse (no QND)
- threshold readout

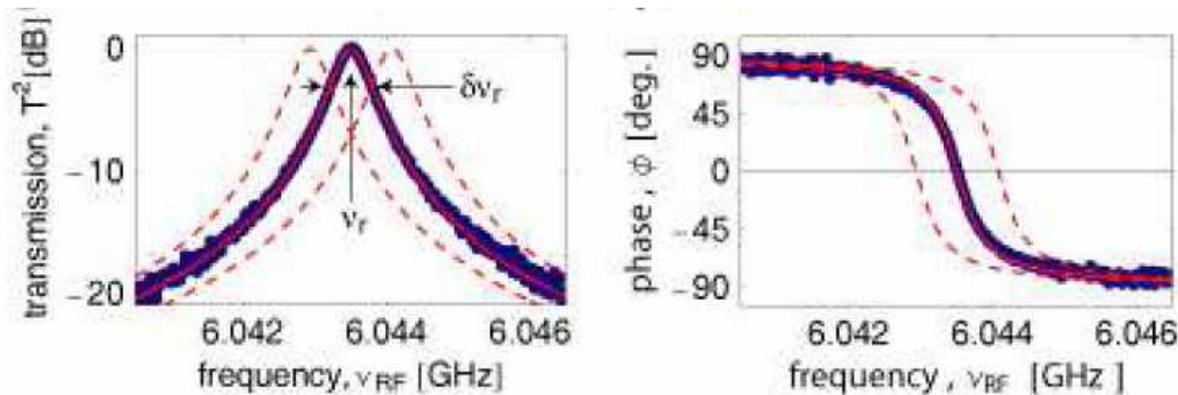
Quantum readout for Josephson qubits

dispersive readout – monitor the eigen-frequency of LC-oscillator



monitor reflection / transmission
amplitude / phase

Sillanpää et al. 2005



Wallraff et al. 2004

Quantum readout for Josephson qubits

Josephson bifurcation amplifier (Siddiqi et al.) – dynamical switching

exploits **nonlinearity** of JJ

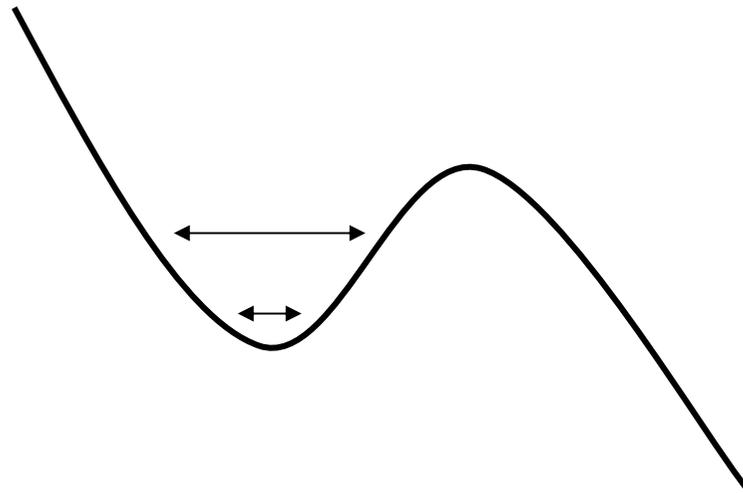
switching between two oscillating states
(different amplitudes and phases)

advantages:

no dc voltage generated, close to QND

higher repetition rate

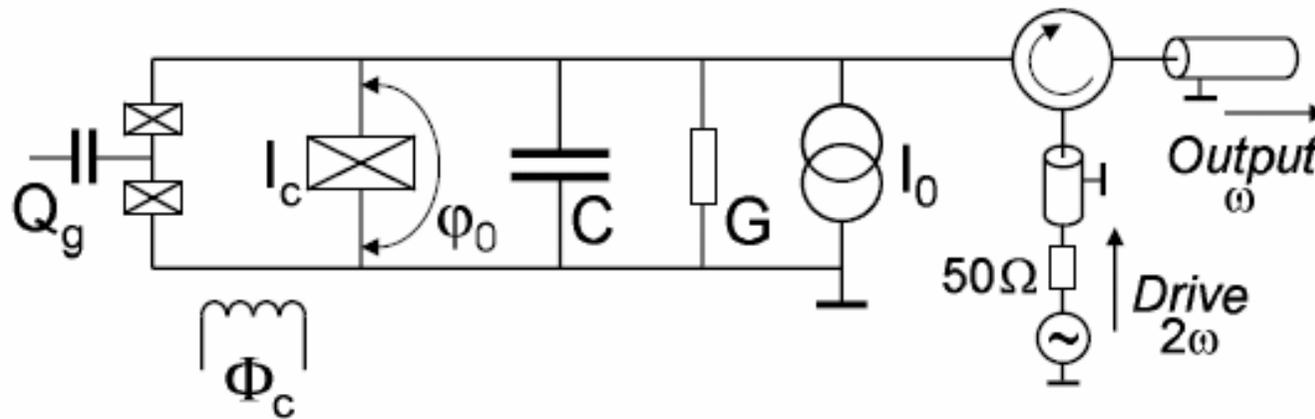
qubit always close to optimal point



Siddiqi et al. 2003

cf. Ithier, thesis 2005

Parametric bifurcation for qubit readout



$$\frac{\hbar C}{2e} \frac{d^2 \varphi}{dt^2} + \frac{\hbar G}{2e} \frac{d\varphi}{dt} + I_c \sin \varphi = I_0 + I_{ac}$$

$$I_0 = I_c \sin \varphi_0$$

$$\varphi = \varphi_0 + x$$

$$\tau = \omega t$$

$$\ddot{x} + x = \xi x - 2\theta \dot{x} + \beta x^2 + \gamma x^3 - \mu x^4 + 3P \cos 2\tau.$$

$$\xi = 1 - \kappa \ll 1, \quad \theta = G/2\omega C \equiv 1/2Q \ll 1,$$

$$\kappa = (\omega_p/\omega)^2$$

$$\beta = 12\mu = (\kappa \tan \varphi_0)/2, \quad \gamma = \kappa/6, \quad 3P = \kappa I_A/I_c$$

Landau, Lifshitz, Mechanics
diff. driving: Dykman et al. '98

Method of slowly-varying amplitudes

$$x \equiv y - P \cos 2\tau$$

$$y = u \cos \tau + v \sin \tau, \quad \dot{y} = -u \sin \tau + v \cos \tau.$$

$$\dot{u} = -\theta u - \frac{1}{2}(\beta P + \xi + \frac{3}{4}\gamma A^2 + \frac{3}{2}\gamma P^2)v - \mu(Pv^3 + \frac{3}{4}P^3v),$$

$$\dot{v} = -\theta v - \frac{1}{2}(\beta P - \xi - \frac{3}{4}\gamma A^2 - \frac{3}{2}\gamma P^2)u - \mu(Pu^3 + \frac{3}{4}P^3u),$$

$$A^2 = u^2 + v^2$$

Equations of motion

$$\begin{aligned}\dot{u} &= -\frac{\partial H}{\partial v} - \theta u + \xi_u(t), \\ \dot{v} &= \frac{\partial H}{\partial u} - \theta v + \xi_v(t).\end{aligned}$$

Hamiltonian

$$H = \frac{\tilde{\xi}}{4}A^2 - \frac{\tilde{\beta}P}{4}A^2 \cos 2\varphi + \frac{3}{32}\gamma A^4 - \frac{\mu P}{4}A^4 \cos 2\varphi,$$

where $\tilde{\beta} = \beta + \frac{3}{2}\mu P^2$, $\tilde{\xi} = \xi + \frac{3}{2}\gamma P^2$, $A^2 = u^2 + v^2$, $u = A \cos \varphi$, $v = A \sin \varphi$

Equations of motion in polar coordinates:

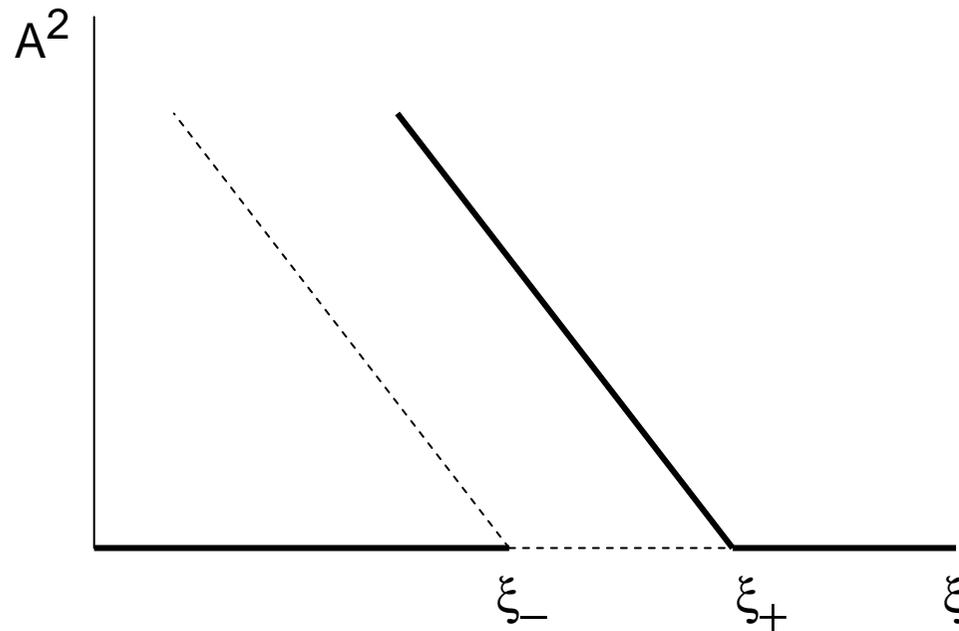
$$\begin{aligned}\frac{d}{dt}(A^2) &= -PA^2 \sin 2\varphi(\tilde{\beta} + \mu A^2) - 2\theta A^2, \\ \dot{\varphi} &= -\frac{\tilde{\beta}P}{2} \cos 2\varphi - \mu PA^2 \cos 2\varphi + \frac{\tilde{\xi}}{2} + \frac{3}{8}\gamma A^2.\end{aligned}$$

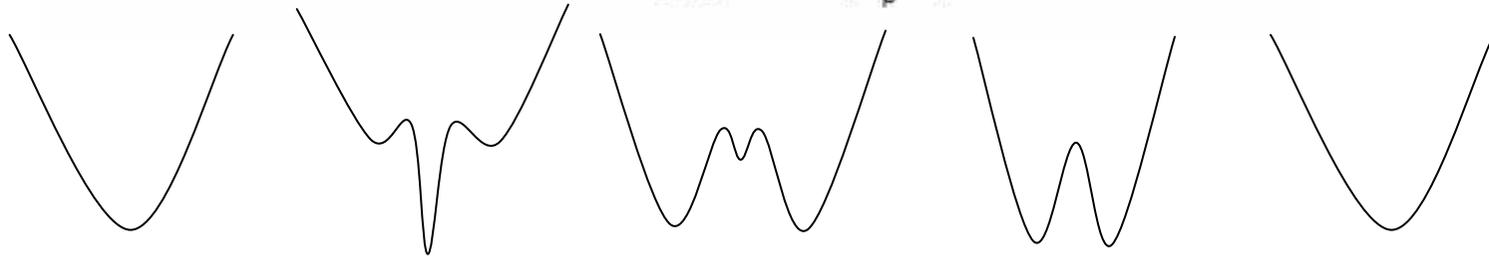
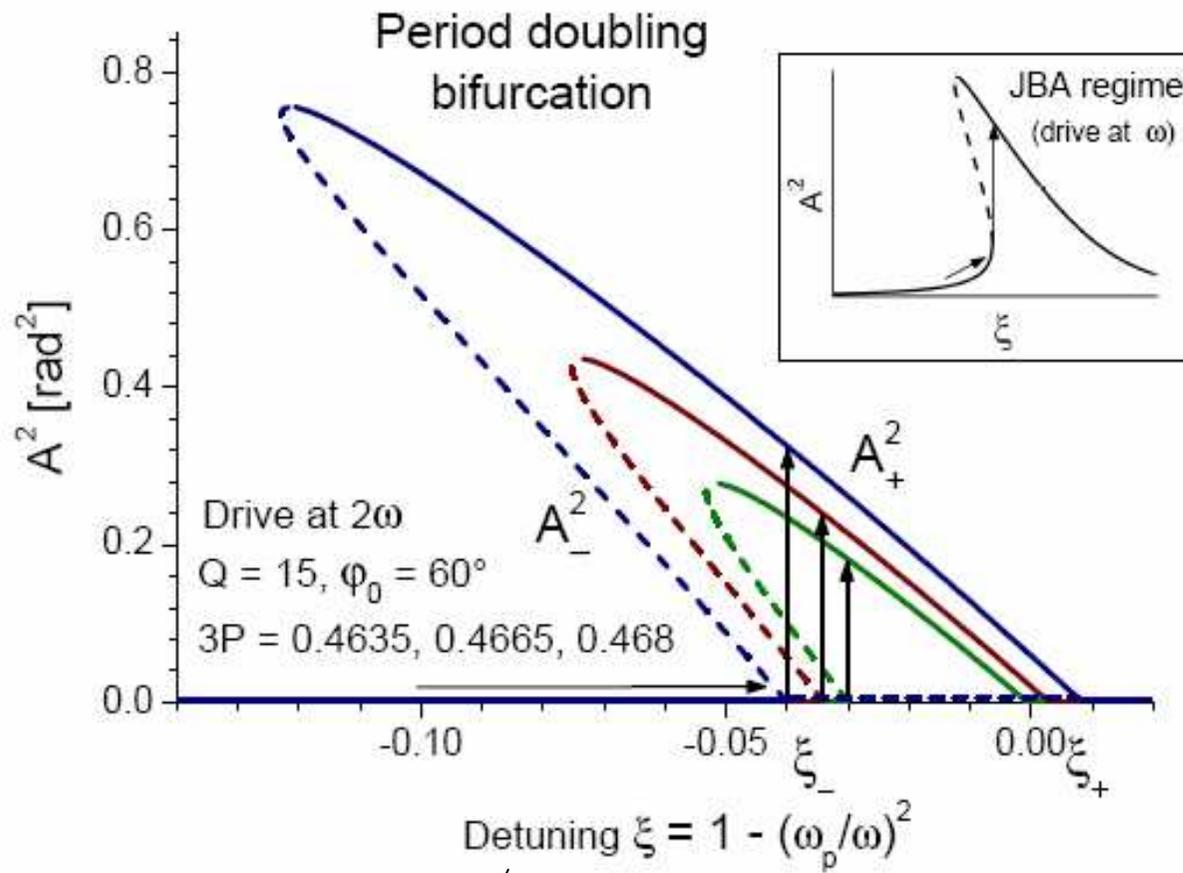
without quartic term

$$\mu=0$$

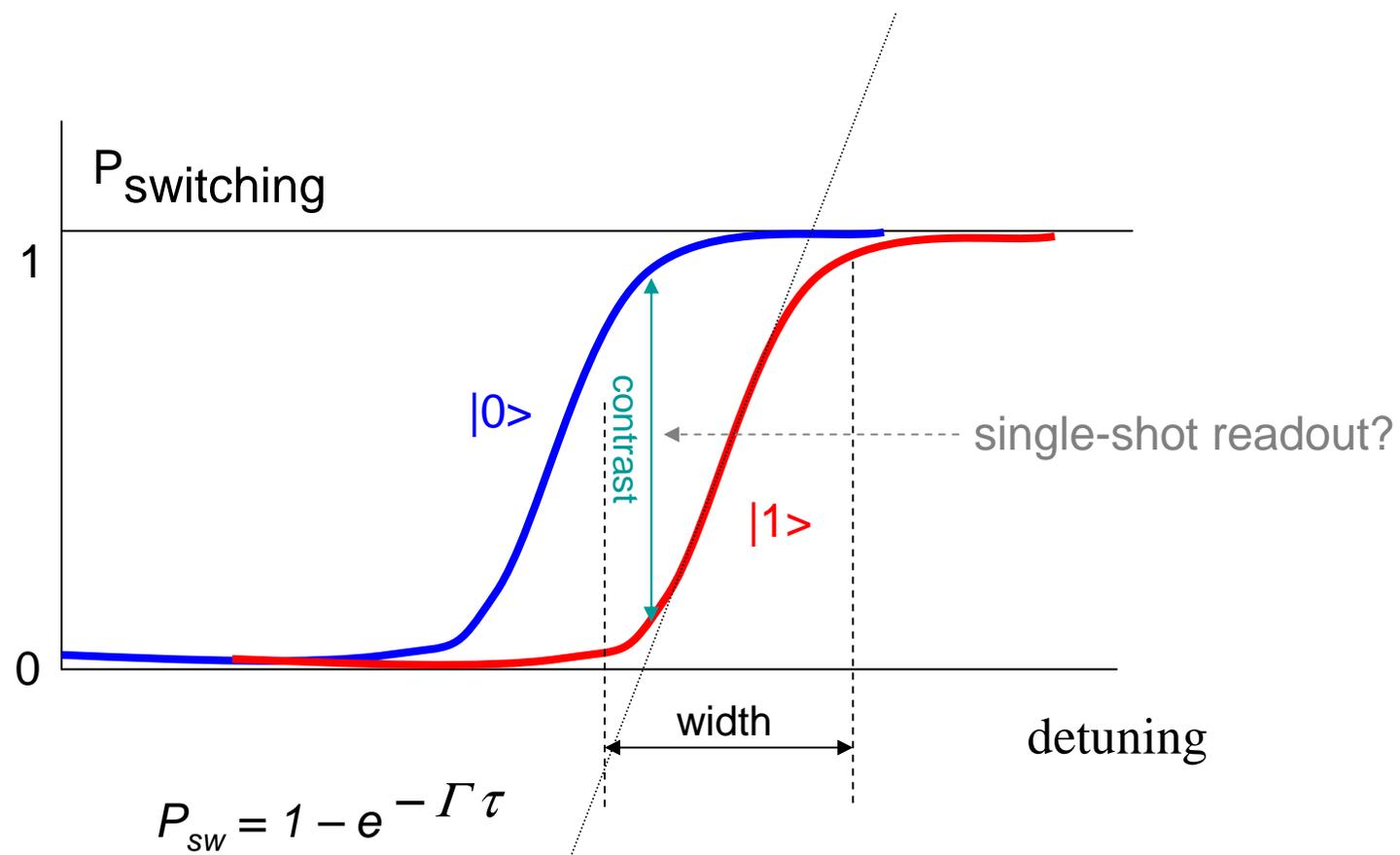
stationary solutions:

$$A=0 \quad \text{or} \quad A_{\pm}^2 = \frac{4}{3\gamma} \left(-\xi - \frac{3}{2}\gamma P^2 \pm \sqrt{\beta^2 P^2 - 4\theta^2} \right)$$

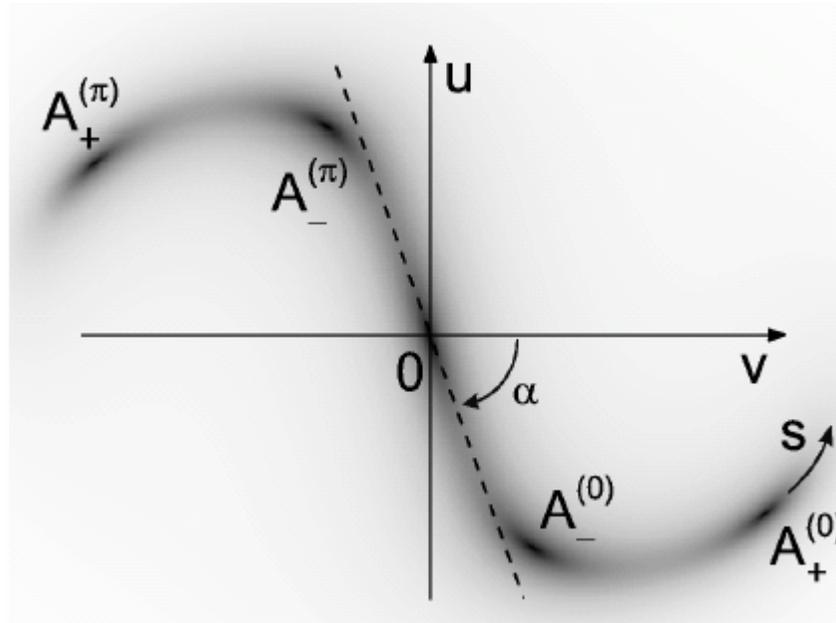




Switching curves



Velocity profile and stationary oscillating states



$0, A_+$ – stable states

A_- – unstable state

near origin: φ relaxes fast, A – slow degree of freedom
(at rate θ)

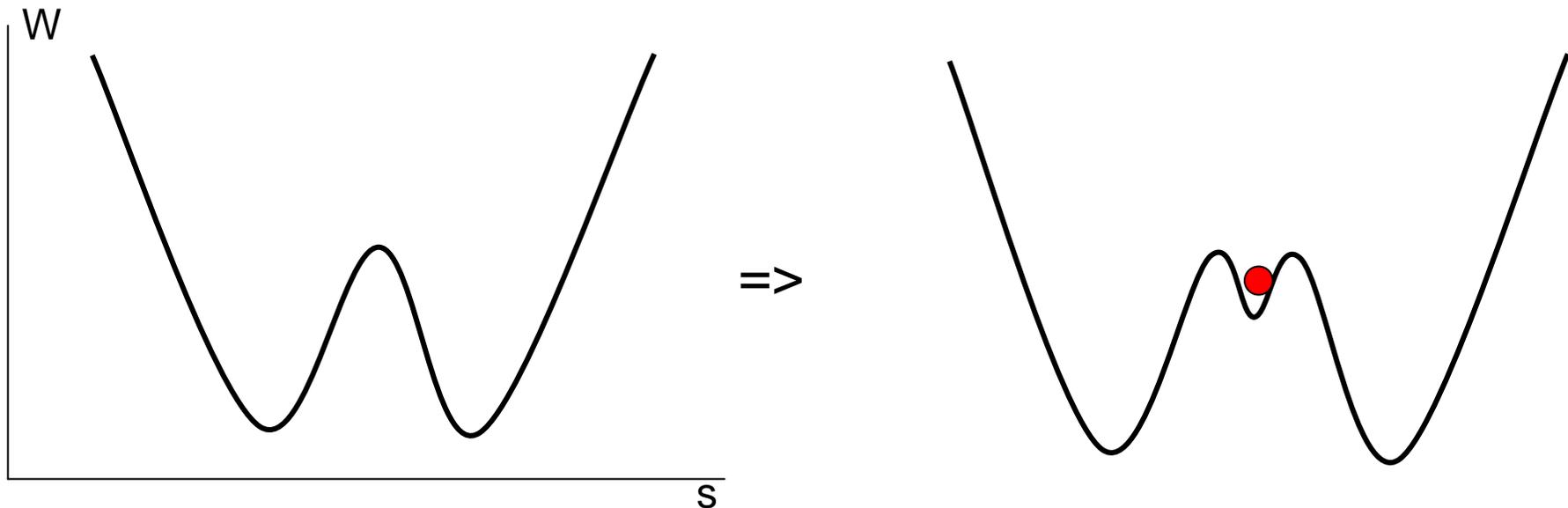
2D Focker-Planck eq. \Rightarrow 1D FPE

Near bifurcation $(\xi \approx \xi_-)$

$$\cos 2\alpha = (\xi + 3/2 \gamma P^2) / P(\beta + 3/2 \mu P^2)$$

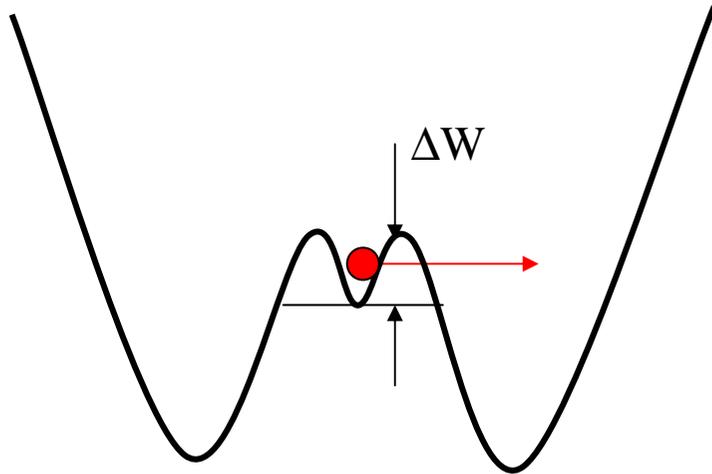
$$ds/dt = -dW(s)/ds + \xi(s)$$

$$W(s) \neq U(s) !$$



Tunneling, switching curves

Focker-Planck equation for P(s)



$$\Gamma \sim \exp(-\Delta W/T_{\text{eff}})$$

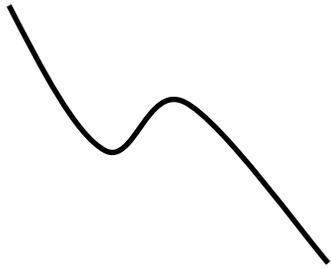
$$W = a s^2 - b s^4$$

$$\Delta W = \frac{1}{12\gamma\theta} \sqrt{(\tilde{\beta}P)^2 - 4\theta^2} (\xi - \xi_-)^2$$

$$T_{\text{eff}} = GT\omega \left(\frac{\omega_p^2}{\omega^2 I_c \cos \varphi_0} \right)^2$$

width of switching curve $\xi - \xi_- \sim \sqrt{T}$

For a generic bifurcation (incl. JBA)



$$U(x) = \varepsilon x - ax^3$$

$$\Delta U \sim \varepsilon^{3/2}$$

controls

$$\Gamma \propto \exp(-\Delta U/T)$$

width of switching curve

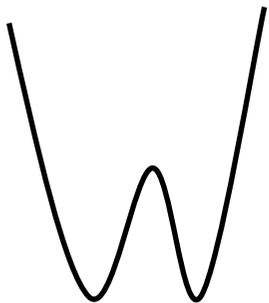
$$\delta\varepsilon \propto T^{2/3}$$

For a parametric bifurcation - additional symmetry

at origin

$$u, v \rightarrow -u, -v$$

shift by period of the drive



$$U = \varepsilon x^2 - ax^4$$

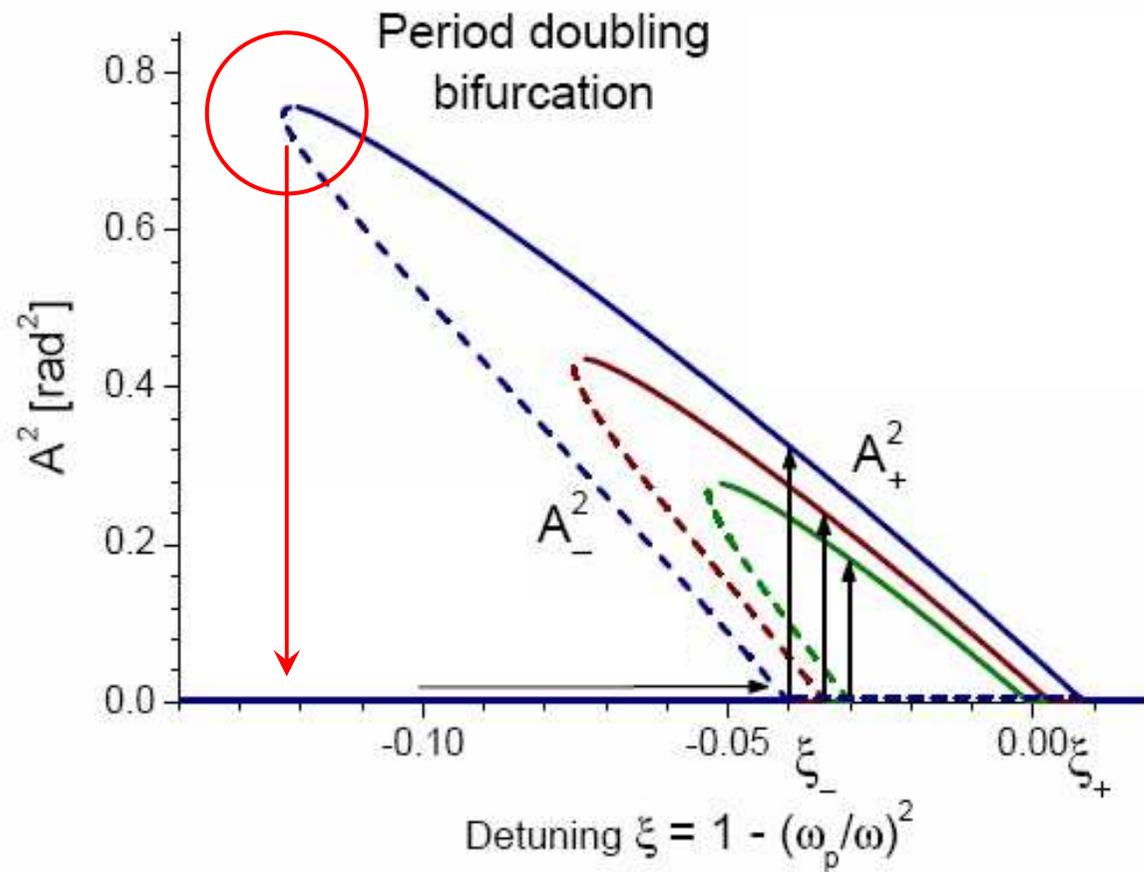
should be even !

$$\Delta U \propto \varepsilon^2$$

=>

$$\delta\varepsilon \propto \sqrt{T}$$

Another operation mode



no mirror symmetry => generic case, stronger effect of cooling

Conclusions

Period-doubling bifurcation readout:

- towards quantum-limited detection?
- low back-action
- rich stability diagram
- various regimes of operation
- tuning amplitude or frequency

results

- bifurcations, tunnel rates,
- switching curves,
- stationary states for various parameters,
- temperature and driving dependence
of the response with double period