

# Emergent chirality

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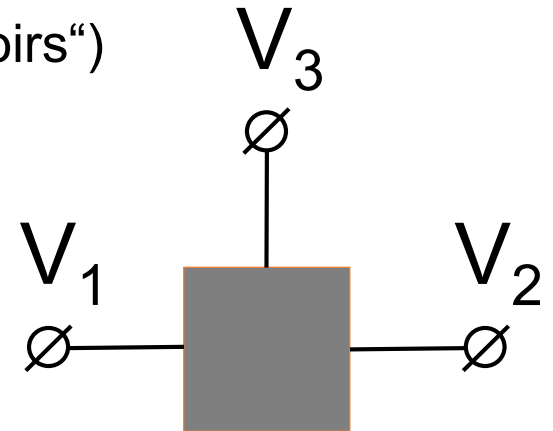
## Outline:

- **emergent chirality on the conceptual level**
- **emergent chirality for a Luttinger-liquid model on the microscopic level**

*largely based on:* D.Aristov, I.Gornyi, D.Polyakov, and P.Wölfle,  
Phys. Rev. B **100**, 165410 (2019)

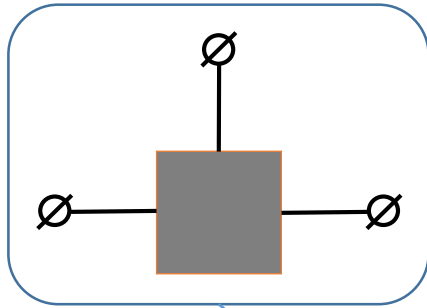
## Main conceptual question for this talk :

- Consider an electron system the **Hamiltonian** of which is **parity** (right-left) and **time-reversal symmetric**
- Connect the system to contacts („ideal reservoirs“) in a right-left symmetric way
- Apply static voltages  $V_1, V_2, V_3 \dots$  to contacts 1, 2, 3 ...



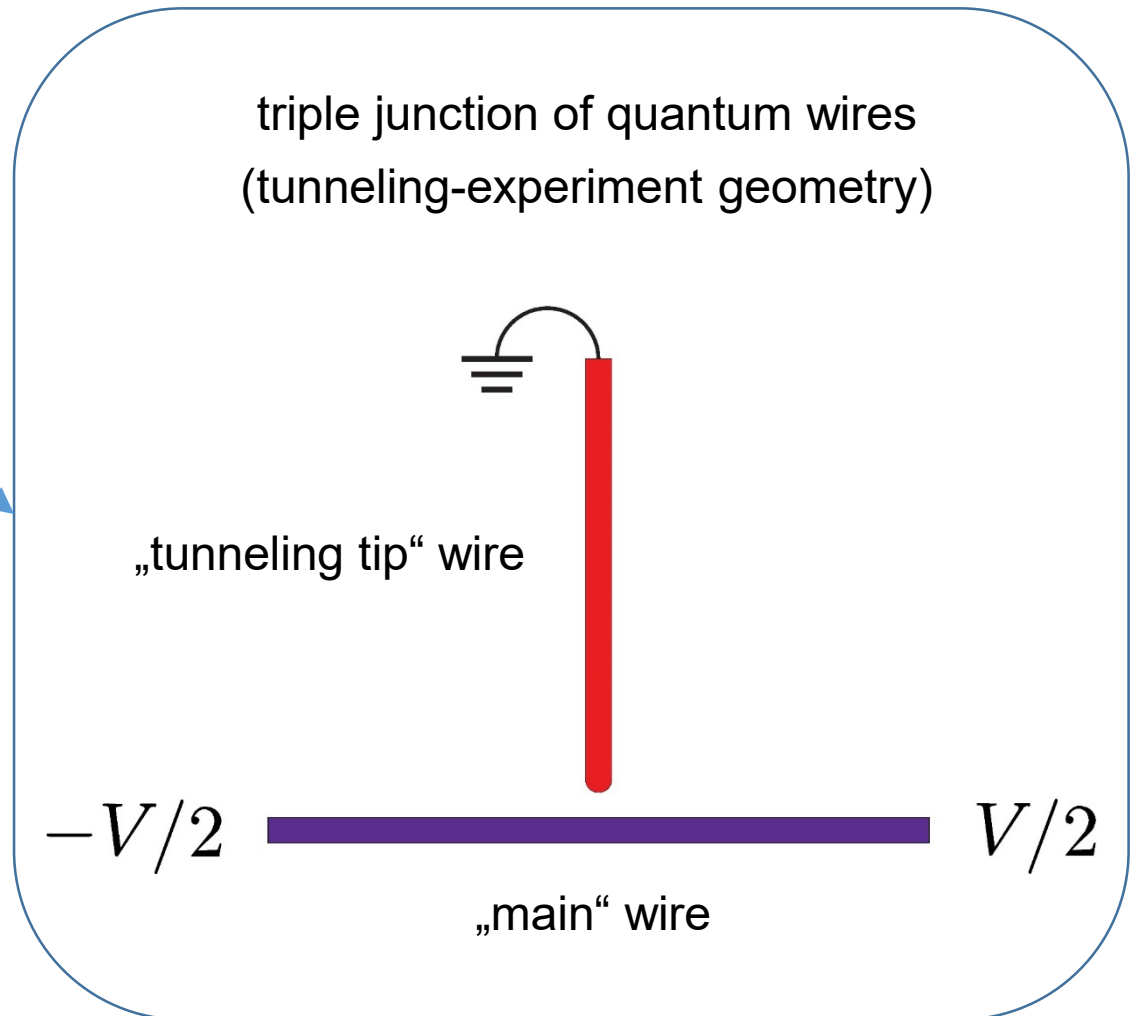
**Can the current response break P and/or T symmetry :**  
**„photogalvanic“ (ratchet) and/or Hall currents ?**

## Distilled example of „emergent chirality“: Y junction



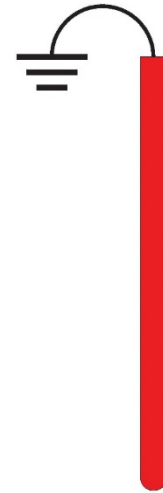
connected  
**symmetrically !**

biased  
**symmetrically !**



## Distilled example of „emergent chirality“: Y junction

zero current in the grounded tip wire ?



connected  
**symmetrically !**

biased  
**symmetrically !**

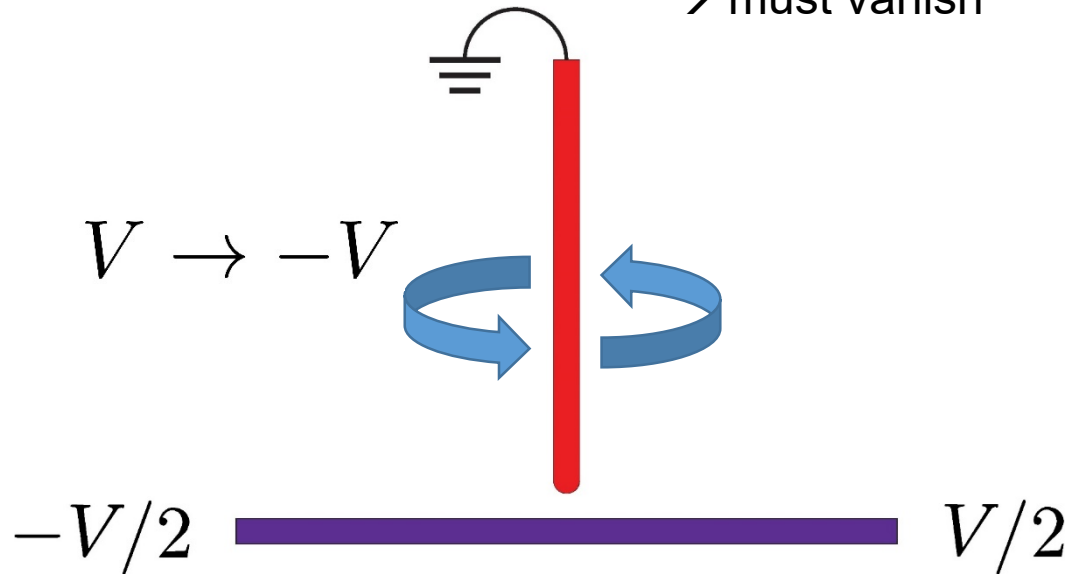
$-V/2$

$V/2$

## Distilled example of „emergent chirality“: Y junction

**zero current in the grounded tip wire ?**

linear response (FDT): current linear in  $V$   
→ must vanish



connected  
**symmetrically !**

biased  
**symmetrically !**

## Distilled example of „emergent chirality“: Y junction

incl.  $e-h$  symmetry  
wrt ground



connected  
**symmetrically !**

biased  
**symmetrically !**

**zero current in the grounded tip wire ?**

exact compensation: for noninteracting electrons

*more current  
to the left*

*more current  
from the right*

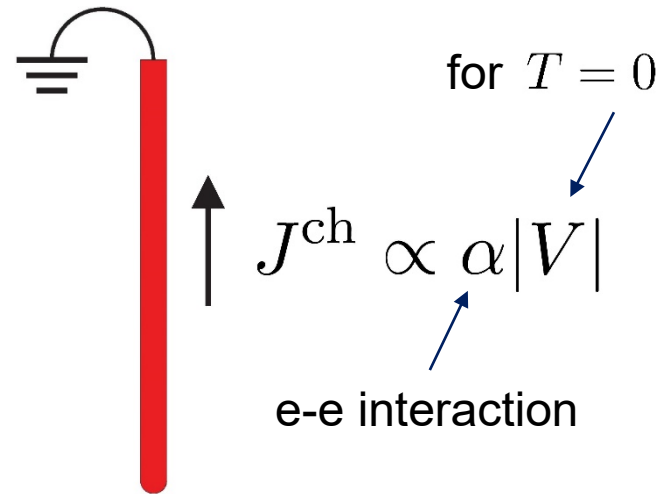
$-V/2$



$V/2$

# Distilled example of „emergent chirality“: Y junction

nonzero current in the tip wire



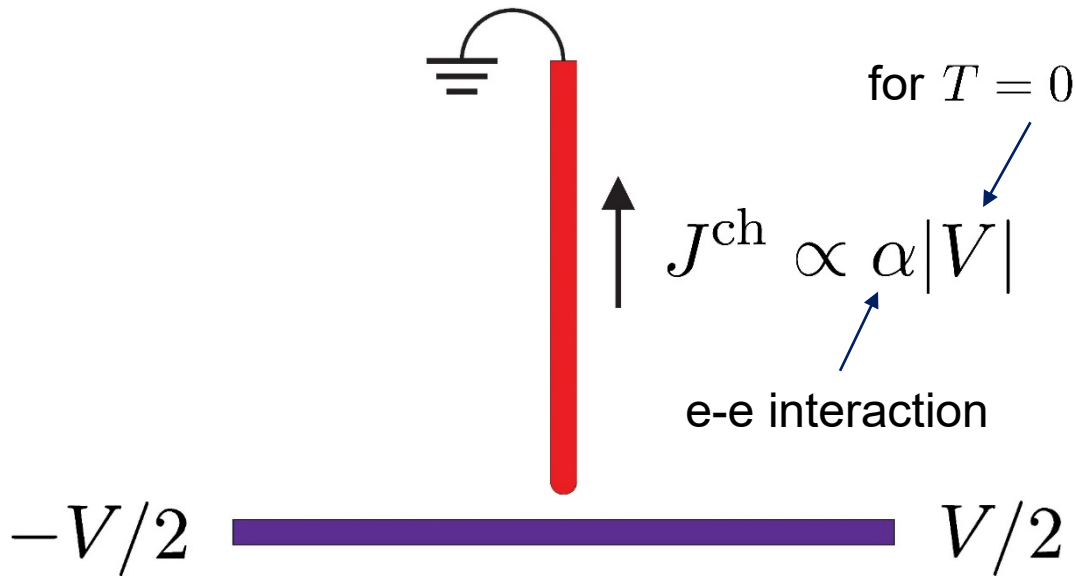
connected  
**symmetrically !**

biased  
**symmetrically !**

$-V/2$   $V/2$

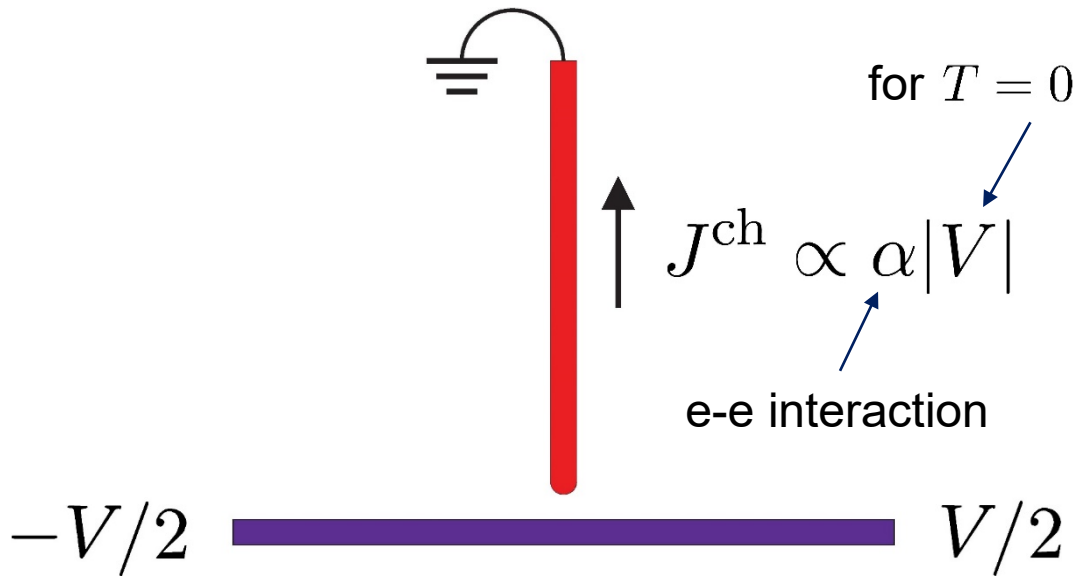


## Chiral current



- **unidirectional** :  $J_{\text{ch}}(V_1, V_2, V_3) = J_{\text{ch}}(-V_1, -V_2, -V_3)$   
≡ chiral flipping **all** voltages
- induced by e-e **interaction** :  $J_{\text{ch}} \propto \alpha$  , depends on the sign  
of interaction
- essentially **nonequilibrium** phenomenon :  $J_{\text{ch}} \propto |V|$

## Chiral current



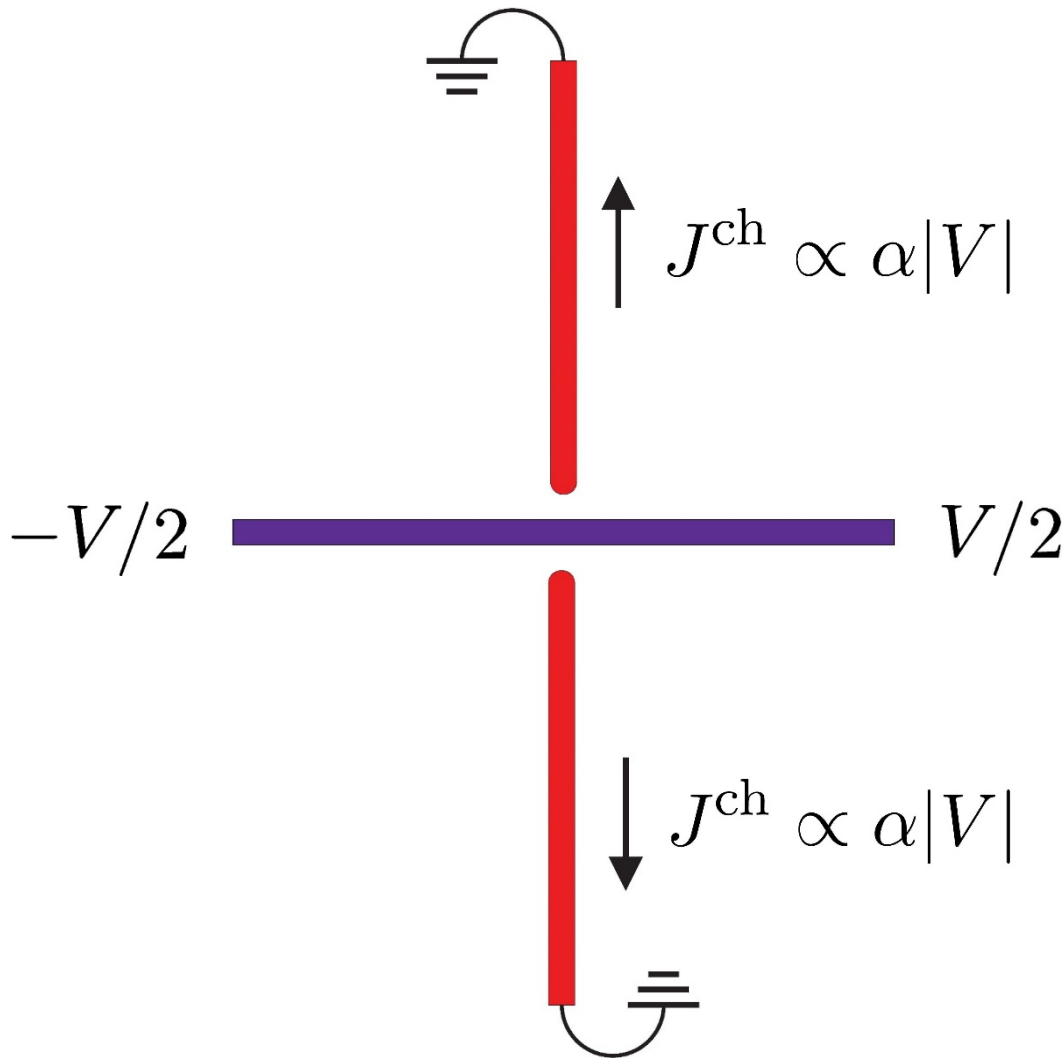
noninteracting tip wire  
( $\alpha$  - in the main wire):

← „**photogalvanic**“  
chiral current

„**Hall**“ chiral current  
→ later

- **unidirectional**:  $J_{\text{ch}}(V_1, V_2, V_3) = J_{\text{ch}}(-V_1, -V_2, -V_3)$   
≡ chiral flipping **all** voltages
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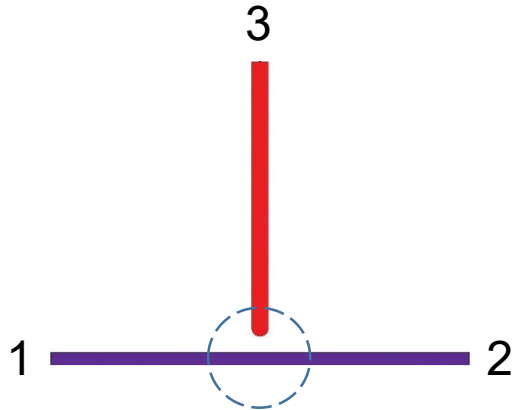
## Chiral current



full „inversion symmetry“

current in the main wire  
ejects ( or absorbs,  
depending on the sign of  $\alpha$  )  
electrons on both „sides“

# Hamiltonian



scattering channels  
= wire indices  
= **contact** indices

Model: **Y junction of ballistic Luttinger liquids**  
(spinless)

$$H_j = \sum_{\eta_j = \pm} \int_0^L dx v_j \left( -i\eta_j \Psi_{\eta_j}^\dagger \partial_x \Psi_{\eta_j} + \pi\alpha_j n_{\eta_j} n_{-\eta_j} \right)$$

wire index

chirality (right- / left-movers)

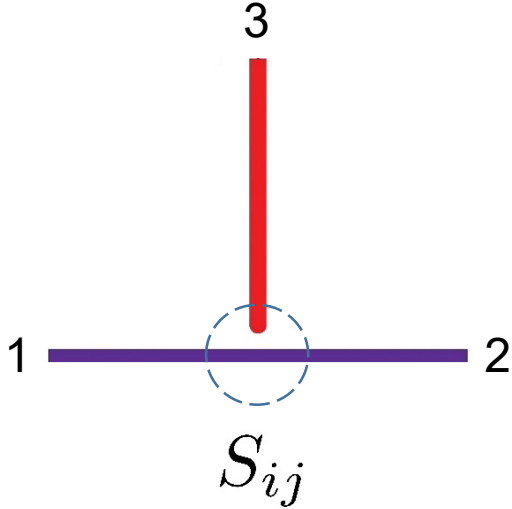
linear dispersion  
(incl. forward e-e scattering)

interaction  
(short-ranged)

- single-particle scattering matrix  $\hat{S}$  at the junction
- boundary condition (on the Keldysh  $G$ ) at the contact:

$$\text{emitted electrons} \rightarrow f_j(\epsilon) = \frac{1}{e^{(\epsilon - \mu_j)/T} + 1}$$

## Noninteracting scattering matrix



scattering channels  
= wire indices  
= **contact** indices

$$\hat{S} = \begin{pmatrix} r & t & t_3 \\ t & r & t_3 \\ t_3 & t_3 & r_3 \end{pmatrix}$$

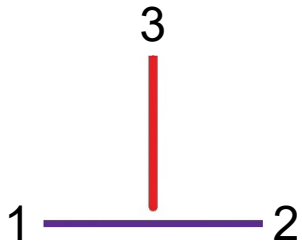
- symmetric (time-reversal symmetric):  $S_{ij} = S_{ji}$
- right-left parity symmetric:  $S_{13} = S_{23}$

unitarity  $\rightarrow$  3 independent parameters  
up to an unobservable global phase

ballistic wires: one of the phases drops out  
from the junction conductances

**noninteracting** Y junction: characterized by **two** conductances

# Scattering off nonequilibrium Friedel oscillations



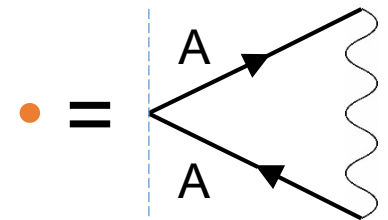
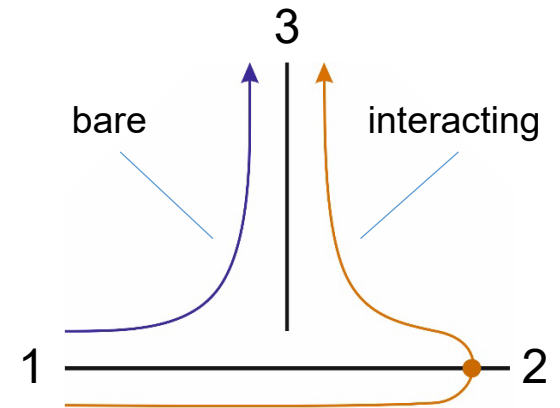
Friedel osc's in wire  $j \leftarrow$  distribution function from contact  $j$  :

$$U_H(x) = \alpha_H 2\text{Re} \int_0^\infty \frac{d\epsilon}{2\pi} e^{-2i\epsilon x/v} \times \begin{cases} r f_1(\epsilon), & x < 0 \\ r^* f_2(\epsilon), & x > 0 \end{cases}$$

Hartree potential in the main (1+2) wire

interfering waves for scattering  $1 \rightarrow 3$  at  $\mathcal{O}(\alpha_H)$ ,  
with interaction in wire 2 :

$$\delta t_3(\epsilon) = \frac{\alpha_H}{2} r^* t t_3 \int_0^\infty d\epsilon' f_2(\epsilon') \frac{1}{\epsilon - \epsilon' + i0}$$



## Mechanism of emergent chirality

- e-e interaction  $\rightarrow$
- **virtual** transitions (screening)
  - **real** transitions

$$\delta t_3^a(\epsilon) \propto \int d\epsilon' f_2(\epsilon') \frac{1}{\epsilon - \epsilon' + i0} \rightarrow \ln \frac{\Lambda}{\max\{T, |\epsilon - \mu_2|\}} - i\pi f_2(\epsilon)$$

„Luttinger liquid renormalization“  
(zero-bias anomaly,  
orthogonality catastrophe, etc.)

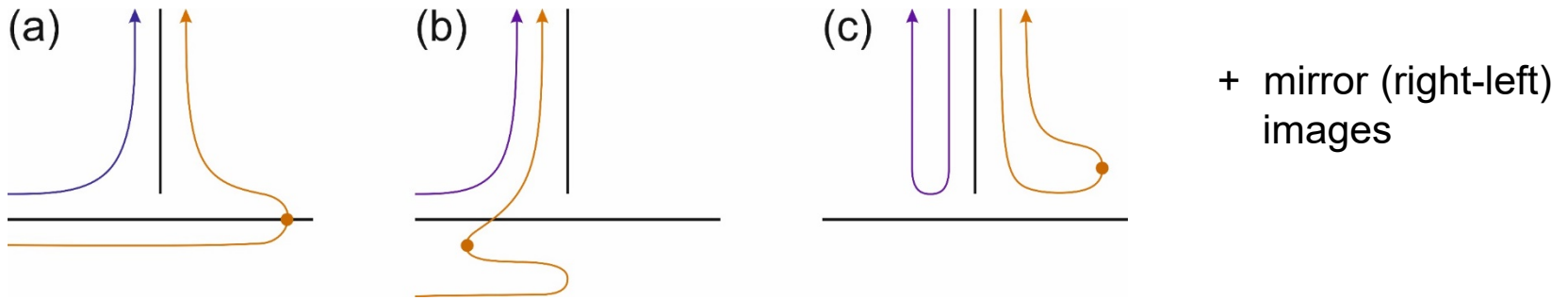
currents **odd** in voltages

**elastic real  
transition**

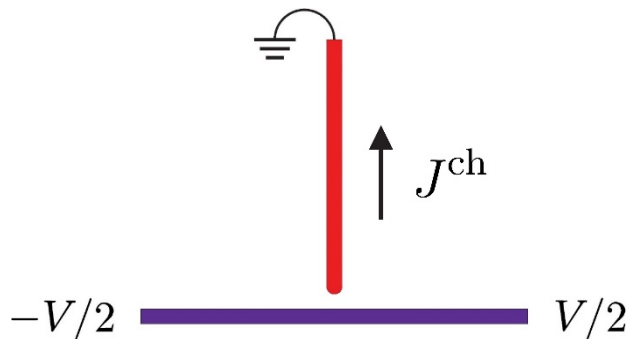
currents **even** in voltages

## Mechanism of emergent chirality

interfering waves for scattering into the tip wire at  $\mathcal{O}(\alpha)$ ,  
with interaction in the main wire :



pole terms in (a)+(b)+(c) for the symmetrically biased junction :  $J^{\text{ch}} = -\frac{e\alpha|V|}{4}\mathcal{A}$



„chirality coefficient“  $\mathcal{A} = \text{Im} \{ |t_3|^2 t r^* \}$

$$T \gg |V| : J^{\text{ch}} \simeq -\frac{1}{8}e\alpha \mathcal{A} \frac{V^2}{T}$$

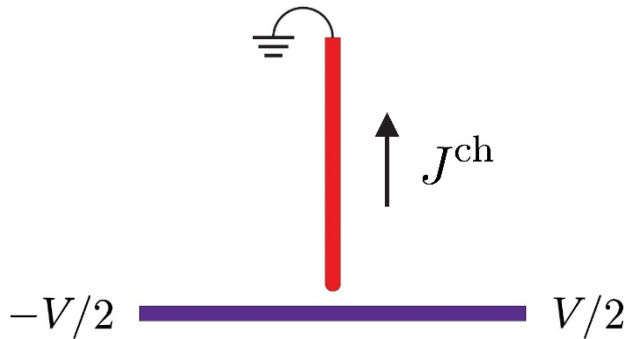


## Mechanism of emergent chirality

chirality coefficient  $\mathcal{A} = \text{Im} \{ |t_3|^2 t r^* \} = \frac{1}{4} \sin^2 \theta \cos \theta \sin \psi$

generically nonzero ■

for a junction of  $N$  (half-)wires with  $N \geq 3$



Y junction  $\rightarrow \mathcal{A} = 0$  :

- decoupled main wire
- perfectly absorbing main wire
- pointlike coupling between the wires

## „Fundamental“ conductance matrix and emergent chirality

$\forall$  (not necessarily symmetric) voltages :  $G_{jk} = e \partial J_j / \partial \mu_k$

redundancy :  $\sum_j G_{jk} = 0$  ,  $\sum_k G_{jk} = 0$  ( Kirchhoff's laws )

$3 \times 3 \hat{G} \rightarrow \text{rank } 2 \rightarrow$  generically **four** independent conductances

$\rightarrow 2 \times 2$  „fundamental“ conductance matrix  $\hat{\tilde{G}} = \begin{pmatrix} G_a & G_c + G_d \\ -G_c + G_d & G_b \end{pmatrix}$

in the  $(a, b)$  basis of independent currents/voltages :

interactions change the structure of  $\hat{\tilde{G}}$  :

off-diagonal  $G_c = G_d = 0$   
for P and T symmetric  
noninteracting electrons

$$J_a = (J_1 - J_2)/2$$

$$J_b = -J_3$$

$$V_a = \mu_1 - \mu_2$$

$$V_b = (\mu_1 + \mu_2)/2 - \mu_3$$

## „Fundamental“ conductance matrix and emergent chirality

$$\hat{\tilde{G}} = \begin{pmatrix} G_a & G_c + G_d \\ -G_c + G_d & G_b \end{pmatrix}$$

**off-diagonal elements**  $\leftarrow$  induced by e-e interaction  
**break P  $\vee$  T symmetry**

$$G_d = \frac{1}{2}(G_{11} - G_{22}) \quad \text{breaks P symmetry} \\ \text{(side diversion aka „photogalvanic“ response)}$$

$$G_c = \frac{1}{2}(G_{12} - G_{21}) \quad \text{breaks T symmetry (Hall response)}$$

def: emergent chirality  $\equiv G_c \vee G_d$  nonzero **and** chiral

$$G_d^{\text{ch}} = \frac{1}{8}e^2 \mathcal{A} [(\alpha_1 + \alpha_2)s_{12} - \alpha_3(s_{23} + s_{31})]$$

$$G_c = -\frac{1}{8}e^2 \mathcal{A} [\alpha_1(s_{12} + s_{31}) + \alpha_2(s_{12} + s_{23}) + \alpha_3(s_{23} + s_{31})]$$

$$s_{jk} = \text{sgn}(\mu_j - \mu_k)$$

## Higher-order renormalization

**chiral** conductances: RG near the fixed points  
at symmetric bias of the main wire ( $V_b = 0$ )

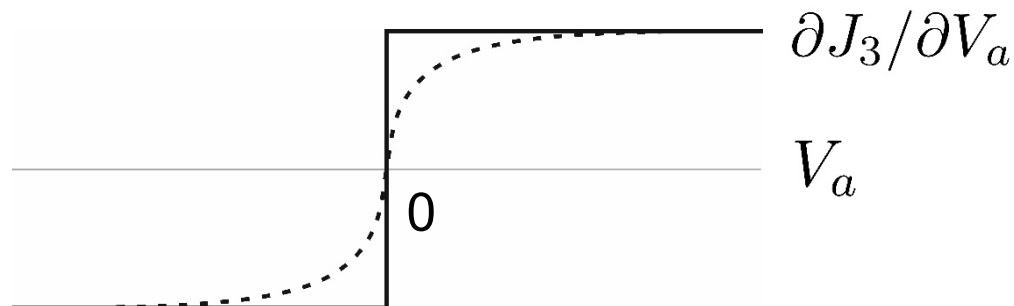
stable fixed point **N** (all wires decoupled):  $J_3 \propto |V_a|^{1+2\alpha+\alpha_3}$

unstable fixed point **A** (tunneling tip decoupled):  $J_3 \propto |V_a|^{1-\alpha+\alpha_3}$

**N** : steplike jump (at  $T=0$ ) of the differential chiral conductance

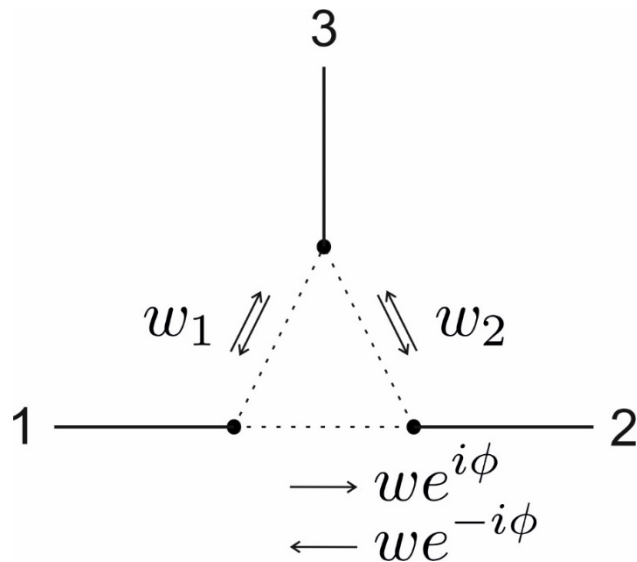
$$e \partial J_3 / \partial V_a = G_c - G_d$$

„smoothed“ by renormalization for repulsive interaction



## Effective chiral model

Y junction  $\rightarrow$  three end points (of wires 1,2,3) connected by hopping



Y junction pierced by a magnetic flux  $\phi$

(model to study the interplay of flux and e-e interaction, attractive interaction-induced „ideal circulator“)

*Oshikawa, Chamon, Affleck '06*

consider the **noninteracting** model **with flux** that mimics (same conductance matrix) the **interacting** model **without flux**

$$\hat{S} = (1 - i\hat{W})^{-1}(1 + i\hat{W})$$

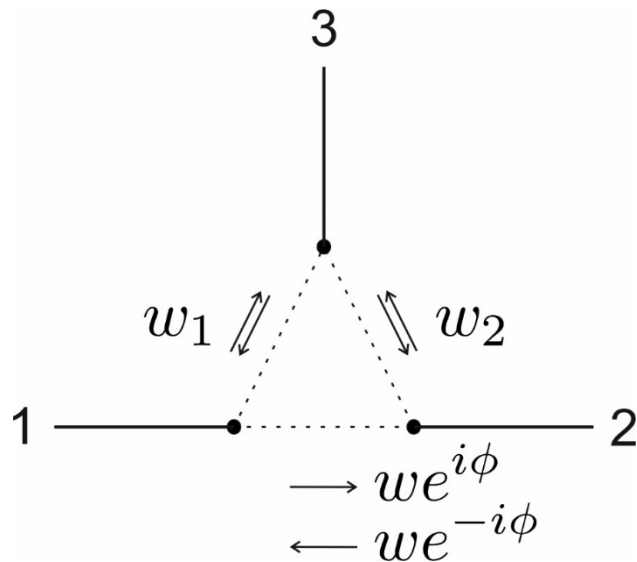
matrix of hopping amplitudes  $\hat{W} = \begin{pmatrix} 0 & we^{-i\phi} & w_1 \\ we^{i\phi} & 0 & w_2 \\ w_1 & w_2 & 0 \end{pmatrix}$

## Effective chiral model

Y junction pierced by a magnetic flux  $\phi$

let  $w_1 = w_2 = w$

hopping amplitudes  $\hat{W} = w \begin{pmatrix} 0 & e^{-i\phi} & 1 \\ e^{i\phi} & 0 & 1 \\ 1 & 1 & 0 \end{pmatrix}$



**eff. flux**  $\leftarrow$  interaction strength,  
changes sign when flipping the sign  
of voltage between terminals 1 and 2

near point **N** :

$$\phi \rightarrow \frac{\pi}{2} \alpha_3 s$$

$$s = \text{sgn} V_a \times s \times \text{sgn} w$$

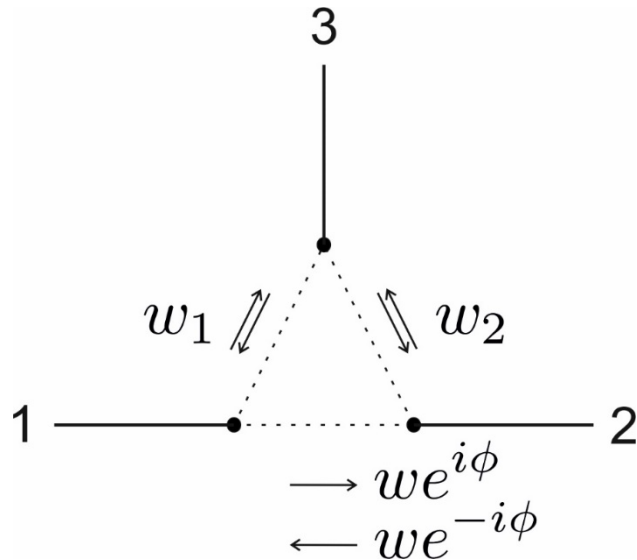
$$s = \text{sgn}(\cos \theta \sin \psi)$$

bare scattering phases

## Effective chiral model

Y junction with built-in asymmetry between  
wires 1 and 2

let  $w_1 \neq w_2$



**eff. anisotropy**  $\leftarrow$  interaction strength,  
changes sign when flipping the sign  
of voltage between terminals 1 and 2

$$\frac{w_1 - w_2}{w_1 + w_2} \rightarrow -\frac{s\pi}{4} (\alpha + \alpha_3) |t| \operatorname{sgn} V_a$$

$$s = \operatorname{sgn} V_a \times s \times \operatorname{sgn} w$$

$$s = \operatorname{sgn}(\cos \theta \sin \psi)$$

## Relation to other works on „rectification“

- currents in a Y junction that are even wrt one voltage, but change sign when all voltages flip sign → not the chiral currents

$$\mathbf{J}^{\text{ch}}(\{\mathbf{V}\}) = \mathbf{J}^{\text{ch}}(-\{\mathbf{V}\})$$

e.g., Wang, Feldman '11

- Bernoulli effect in fluid-filled pipes

*Bernoulli 1738, Euler 1752, Venturi 1797*

$$\rho v^2 / 2 + p = \text{const} \quad (\text{no gravity})$$

conservation of energy in isentropic incompressible fluid  
 $v$  hydrodynamic velocity,  $\rho$  mass density,  $p$  pressure

fluid is „sucked in“ irrespective of the direction of flow in the „main pipe“

requires fast equilibration in the moving frame (hydrodynamics)  
at each point of the junction

- quantum dot capacitively coupled to the leads + inelastic cotunneling

*Jordan, Büttiker '08*

requires charge quantization on the junction



rehau.com



## Conclusion:

- *emergent chirality*  
*in a system with  $P$  and  $T$  symmetric Hamiltonian :*  
**emergence of currents that break  $P$  and/or  $T$  symmetry  
and do not change when all voltages flip signs**
- *emergent chirality*
  - **induced by interaction between particles**
  - **essentially nonequilibrium phenomenon**