## Canonical conservation laws and integrability conditions for difference equations

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Landau Conference, 21-23 June 2010

- Difference equation Q = 0 and Dynamical variables.
- Difference fields  $\mathcal{F}_Q, \mathcal{F}_0, \mathcal{F}_s, \mathcal{F}_t$ , the elimination map.
- Symmetries and conservation laws
- Recursion operators for difference equations
- Formal difference series, difference Adler Theorem.
- Canonical conservation laws: integrability conditions

Difference equations on  $\mathbb{Z}^2$  can be seen as a discrete analogue of partial differential equations with two independent variables.

Let us denote by u = u(n, m) a complex-valued function  $u : \mathbb{Z}^2 \mapsto \mathbb{C}$  where n and m are "independent variables" and u will play the rôle of a "dependent" variable in a difference equation.

Instead of partial derivatives we have two commuting shift maps  ${\cal S}$  and  ${\cal T}$  defined as

 $\mathcal{S}: u \mapsto u_{1,0} = u(n+1,m), \qquad \mathcal{T}: u \mapsto u_{0,1} = u(n,m+1)$ 

For uniformity of notations, we denote u as  $u_{0,0}$ .

In the theory of difference equations we shall treat symbols  $u_{p,q}$  as variables.

We denote  $U = \{u_{p,q} \mid (p,q) \in \mathbb{Z}^2\}.$ 

For a function 
$$f = f(u_{p_1,q_1}, ..., u_{p_k,q_k})$$
  
 $S^i \mathcal{T}^j(f) = f_{i,j} = f(u_{p_1+i,q_1+j}, ..., u_{p_k+i,q_k+j}).$ 

A quadrilateral difference equation can be defined as

$$Q(u_{0,0}, u_{1,0}, u_{0,1}, u_{1,1}) = 0$$
,

where  $Q(u_{0,0}, u_{1,0}, u_{0,1}, u_{1,1})$  is an **irreducible polyno**mial of the "dependent variable"  $u = u_{0,0}$  and its shifts.

$$Q_{p,q} = Q(u_{p,q}, u_{p+1,q}, u_{p,q+1}, u_{p+1,q+1}) = 0, \qquad (p, q) \in \mathbb{Z}^2$$

We shall assume that Q is an irreducible **affine-linear polynomial** which depends non-trivially on all variables:

$$\frac{\partial Q}{\partial u_{i,j}} \neq 0, \quad \frac{\partial^2 Q}{\partial^2 u_{i,j}} = 0, \qquad i, j \in \{0, 1\}.$$

Let  $\mathbb{C}[U]$  be the ring of polynomials.

 $S, T \in Aut \mathbb{C}[U]$  and thus  $\mathbb{C}[U]$  is a **difference ring**.

The difference ideal

$$J_Q = \langle \{Q_{p,q} \,|\, (p, q) \in \mathbb{Z}^2 \} \rangle$$

is prime and thus the quotient ring  $\mathbb{C}[U]/J_Q$  is an integral domain.

Solutions of the difference equation are points of the affine variety  $V(J_Q)$ .

Rational functions of variables  $u_{p,q}$  on  $V(J_Q)$  form a field

$$\mathcal{F}_Q = \{ [a]/[b] \mid a, b \in \mathbb{C}[U], \ b \notin J_Q \},\$$

where [a] denotes the class of equivalent polynomials (two polynomials  $f, g \in \mathbb{C}[U]$  are equivalent, denoted by  $f \equiv g$ , if  $f - g \in J_Q$ ).

For  $a, b, c, d \in \mathbb{C}[U]$ ,  $b, d \notin J_Q$ , rational functions a/b and c/d represent the same element of  $\mathcal{F}_Q$  if  $ad - bc \in J_Q$ .

The fields of rational functions of variables

 $U_{\mathbf{s}} = \{u_{n,0} \mid n \in \mathbb{Z}\}, \quad U_{\mathbf{t}} = \{u_{0,n} \mid n \in \mathbb{Z}\}, \quad U_{\mathbf{0}} = U_{\mathbf{s}} \cup U_{\mathbf{t}}.$ are denoted respectively as

$$\mathcal{F}_{s} = \mathbb{C}(U_{s}), \qquad \mathcal{F}_{t} = \mathbb{C}(U_{t}), \qquad \mathcal{F}_{0} = \mathbb{C}(U_{0}).$$

In the affine-linear case we can uniquely resolve equation Q = 0 with respect to each variable

$$\begin{split} u_{0,0} &= F(u_{1,0}, u_{0,1}, u_{1,1}), & u_{1,0} = G(u_{0,0}, u_{0,1}, u_{1,1}), \\ u_{0,1} &= H(u_{0,0}, u_{1,0}, u_{1,1}), & u_{1,1} = M(u_{0,0}, u_{1,0}, u_{0,1}). \end{split}$$
  
We can recursively and uniquely express any variable  $u_{p,q}$  in terms of the variables  $U_0 = U_{\mathbf{s}} \cup U_{\mathbf{t}}.$ 

**Definition 1.** For elements of U the elimination map  $\mathcal{E}: U \mapsto \mathbb{C}(U_0)$  is defined recursively:

 $\begin{aligned} \forall p \in \mathbb{Z}, & \mathcal{E}(u_{0,p}) = u_{0,p}, & \mathcal{E}(u_{p,0}) = u_{p,0}, \\ \text{if } p > 0, q > 0, & \mathcal{E}(u_{p,q}) = M(\mathcal{E}(u_{p-1,q-1}), \mathcal{E}(u_{p,q-1}), \mathcal{E}(u_{p-1,q})), \\ \text{if } p < 0, q > 0, & \mathcal{E}(u_{p,q}) = H(\mathcal{E}(u_{p,q-1}), \mathcal{E}(u_{p+1,q-1}), \mathcal{E}(u_{p+1,q})), \\ \text{if } p > 0, q < 0, & \mathcal{E}(u_{p,q}) = G(\mathcal{E}(u_{p-1,q}), \mathcal{E}(u_{p-1,q+1}), \mathcal{E}(u_{p,q+1})), \\ \text{if } p < 0, q < 0, & \mathcal{E}(u_{p,q}) = F(\mathcal{E}(u_{p+1,q}), \mathcal{E}(u_{p,q+1}), \mathcal{E}(u_{p+1,q+1})). \end{aligned}$ 

For polynomials  $f(u_{p_1,q_1},\ldots,u_{p_k,q_k}) \in \mathbb{C}[U]$  the elimination map  $\mathcal{E} : \mathbb{C}[U] \mapsto \mathbb{C}(U_0)$  is defined as

 $\mathcal{E}$ :  $f(u_{p_1,q_1}, \ldots, u_{p_k,q_k}) \mapsto f(\mathcal{E}(u_{p_1,q_1}), \ldots, \mathcal{E}(u_{p_k,q_k})) \in \mathbb{C}(U_0)$ . For rational functions a/b,  $a, b \in \mathbb{C}[U]$ ,  $b \notin J_Q$  the elimination map  $\mathcal{E}$  is defined as

$$\mathcal{E}: a/b \mapsto \mathcal{E}(a)/\mathcal{E}(b).$$

Variables  $U_0$  we shall call the **dynamical** variables.

 $\mathcal{E}(u_{p,q})$  is a rational function of |p| + |q| + 1 dynamical variables  $(pq \neq 0)$ 

 $\mathcal{E}(u_{p,q}) \in \mathbb{C}(\{u_{n,0}, u_{0,m} \mid 0 \le |n-p| \le |p|, \ 0 \le |m-q| \le |q|\}).$ 

The elimination map  $\mathcal{E} : \mathbb{C}[U] \mapsto \mathbb{C}(U_0)$  is a **difference** ring homomorphism

$$\operatorname{Ker} \mathcal{E} = J_Q, \qquad \operatorname{Im} \mathcal{E} \sim \mathbb{C}[U]/J_Q.$$

The field  $\mathbb{C}(U_0)$  is a **difference field** with automorphisms  $\mathcal{E} \circ \mathcal{S}$  and  $\mathcal{E} \circ \mathcal{T}$ . The map  $\mathcal{E} : \mathcal{F}_Q \mapsto \mathbb{C}(U_0)$  is a **difference field isomorphism**.

Map  $\mathcal{E}$  is a useful tool to establish whether two rational functions f, g of variables U are equivalent (i.e. represent the same element of  $\mathcal{F}_Q$ ):  $f \equiv g \Leftrightarrow \mathcal{E}(f) = \mathcal{E}(g)$ .

## Symmetries and conservation laws of difference equations

**Definition 2.** Let Q = 0 be a difference equation. Then  $K \in \mathcal{F}_Q$  is called a **symmetry** (a generator of an infinitesimal symmetry) of the difference equation if

 $D_Q(K) \equiv 0.$ 

Here  $D_Q$  is the Frechét derivative of Q defined as

$$D_Q = \sum_{i,j} Q_{u_{i,j}} \mathcal{S}^i \mathcal{T}^j , \qquad Q_{u_{i,j}} = \frac{\partial Q}{\partial u_{i,j}}$$

What one has to check is that  $\mathcal{E}(D_Q(K)) = 0$ .

If K is a symmetry and u = u(n, m) is a solution of a difference equation Q = 0, then the infinitesimal transformation of solution u:

$$\hat{u} = u + \epsilon K$$

satisfies equation

$$Q(\hat{u}_{0,0}, \hat{u}_{1,0}, \hat{u}_{0,1}, \hat{u}_{1,1}) \equiv \mathcal{O}(\epsilon^2).$$

If the difference equation Q = 0 admits symmetries, then they form a Lie algebra. With a symmetry  $K \in \mathcal{F}_Q$  we associate an evolutionary derivation ( $\mathcal{S}X_K = X_K \mathcal{S}, \ \mathcal{T}X_K = X_K \mathcal{T}$ ) of  $\mathcal{F}_Q$  (or a vector field on  $\mathcal{F}_Q$ ):

$$X_K = \sum_{(p,q)\in\mathbb{Z}^2} K_{p,q} \frac{\partial}{\partial u_{p,q}}, \qquad K_{p,q} = \mathcal{S}^p \mathcal{T}^q(K)$$

For any  $a \in J_Q$  we have  $X_K(a) \in J_Q$  and thus the evolutionary derivation  $X_K$  is defined **correctly** on  $\mathcal{F}_Q$ .

$$X_F X_G - X_G X_F = X_H,$$

where H = [F, G] is also a symmetry, with [F, G] denoting the **Lie bracket** 

$$[F,G] = X_F(G) - X_G(F) = D_G(F) - D_F(G) \in \mathcal{F}_Q.$$

The Lie algebra of symmetries of the difference equation Q = 0 will be denoted as  $\mathfrak{A}_Q$ .

Existence of an infinite dimensional Lie algebra  $\mathfrak{A}_Q$  is a characteristic property of integrable equations and can be taken as a **definition of integrability**.

Symmetry K can be represented by  $\mathcal{E}(K) \in \mathcal{F}_0$ .

It is known, that for a quadrilateral equation a symmetry is a sum of two functions

$$K = K_{\mathbf{s}}(u_{N_1,0},\ldots,u_{N_2,0}) + K_{\mathbf{t}}(u_{0,M_1},\ldots,u_{0,M_2})$$

For the ABS equations any "five point symmetry"  $K = K(u_{0,0}, u_{-1,0}, u_{0,-1}, u_{1,0}, u_{0,1})$  is a **sum of symmetries** 

$$K = K_{s}(u_{-1,0}, u_{0,0}, u_{1,0}) + K_{t}(u_{0,-1}, u_{0,0}, u_{0,1}).$$
  
For  $K_{s} \in \mathcal{F}_{s}$ : ord  $K_{s}(u_{N_{1},0}, \dots, u_{N_{2},0}) = (N_{1}, N_{2}).$   
For  $K_{t} \in \mathcal{F}_{t}$ : ord  $K_{t}(u_{0,M_{1}}, \dots, u_{0,M_{2}}) = (M_{1}, M_{2}).$ 

**Definition 3.** (1) A pair  $(\rho, \sigma) \in \widehat{\mathcal{F}}_Q$  is called a **conser**vation law for the difference equation Q = 0, if

$$(\mathcal{T}-1)(\rho) \equiv (\mathcal{S}-1)(\sigma).$$

Functions  $\rho$  and  $\sigma$  will be referred to as the density and the flux of the conservation law and 1 denotes the identity map.

(2) A conservation law is called **trivial**, if functions  $\rho$  and  $\sigma$  are components of a (difference) gradient of some element  $H \in \mathcal{F}_Q$ , i.e.

$$\rho = (S-1)(H), \qquad \sigma = (T-1)(H).$$

(3) If  $\rho_1 - \rho_2 \in \text{Im}(S-1)$ , then  $\rho_1 \cong \rho_2$ .

Typically conserved densities belong to  $\mathcal{F}_s$  or  $\mathcal{F}_t$ .

Euler's operator gives a criteria to determine whether two elements of  $\mathcal{F}_s$  are equivalent or not.

**Definition 4.** Let  $f \in \mathcal{F}_s$  has order  $(N_1, N_2)$ , then the variational derivative  $\delta_s$  of f is defined as

$$\delta_{\mathbf{s}}(f) = \sum_{k=N_1}^{N_2} \mathcal{S}^{-k} \left( \frac{\partial f}{\partial u_{k,0}} \right).$$

For  $\rho, \varrho \in \mathcal{F}_{s}$ ,  $\rho \cong \varrho \Leftrightarrow \delta_{s}(\rho) = \delta_{s}(\varrho)$ . If  $\rho$  is trivial then  $\delta_{s}(\rho) = 0$ .

The order of a density  $\rho \in \mathcal{F}_s$  is defined as  $\operatorname{ord}_{\delta_s}(\rho) = N_2 - N_1$ , where  $(N_1, N_2) = \operatorname{ord}(\delta_s(\rho))$ .

Recursion operators for difference equations Definition 5. (1) Elements of  $\mathcal{F}_Q[S]$  are called s-difference operators. (2) Elements of  $\mathcal{F}_Q(S)$  are called s-pseudo-difference operators.

Similarly one can define t-difference and t-pseudo-difference operators.

The action of a difference operator  $A \in \mathcal{F}_Q[S]$  on elements of  $\mathcal{F}_Q$  is naturally defined and  $\text{Dom}(A) = \mathcal{F}_Q$ .

The domain of a pseudo-difference operator  $B \in \mathcal{F}_Q(S)$ is defined as  $\text{Dom}(B) = \{a \in \mathcal{F}_Q | B(a) \in \mathcal{F}_Q\}.$ For instance, if  $B = FG^{-1}$  where  $F, G \in \mathcal{F}_Q[S]$ , then Dom(B) = Im G. A recursion operator of a difference equation Q = 0is a pseudo-difference operator  $\Re$  such that

 $\mathfrak{R}: \mathsf{Dom}(\mathfrak{R}) \cap \mathfrak{A}_Q \mapsto \mathfrak{A}_Q,$ 

where  $\mathfrak{A}_Q$  is the linear space of symmetries of this difference equation.

In other words, if the action of  $\mathfrak{R}$  on a symmetry  $K \in \mathcal{F}_Q$ is defined, i.e.  $\mathfrak{R}(K) \in \mathcal{F}_Q$ , then  $\mathfrak{R}(K)$  is a symmetry of the same difference equation. **Theorem 1.** Let  $Q(u_{0,0}, u_{1,0}, u_{0,1}, u_{1,1}) = 0$  be a difference equation.

(i) If there exist two s-pseudo-differential operators  $\mathfrak R$  and  $\mathfrak P$  such that

$$D_Q \circ \mathfrak{R} = \mathfrak{P} \circ D_Q,$$

then  $\Re$  is a recursion operator of the difference equation.

(ii) Relation (1) is valid if and only if

$$\mathcal{T}(\mathfrak{R}) - \mathfrak{R} = [\Phi \circ \mathfrak{R}, \Phi^{-1}] ,$$
  
where  $\Phi = (Q_{u_{1,1}}\mathcal{S} + Q_{u_{0,1}})^{-1} \circ (Q_{u_{1,0}}\mathcal{S} + Q_{u_{0,0}})$ , and the operator  $\mathfrak{P}$  satisfies

$$\mathfrak{P} = (Q_{u_{1,0}}\mathcal{S} + Q_{u_{0,0}}) \circ \mathfrak{R} \circ (Q_{u_{1,0}}\mathcal{S} + Q_{u_{0,0}})^{-1}.$$

**Theorem 2.** Let  $Q(u_{0,0}, u_{1,0}, u_{0,1}, u_{1,1}) = 0$  be a difference equation.

(i) If there exist two t-pseudo-differential operators  $\widehat{\mathfrak{R}}$  and  $\widehat{\mathfrak{P}}$  such that

$$D_Q \circ \hat{\mathfrak{R}} = \hat{\mathfrak{P}} \circ D_Q, \tag{1}$$

then  $\widehat{\mathfrak{R}}$  is a recursion operator of the difference equation.

(ii) Relation (1) is valid if and only if  $S(\hat{\Re}) - \hat{\Re} = [\Psi \circ \hat{\Re}, \Psi^{-1}]$ where  $\Psi = (Q_{u_{1,1}}\mathcal{T} + Q_{u_{1,0}})^{-1} \circ (Q_{u_{0,1}}\mathcal{T} + Q_{u_{0,0}})$ , and the operator  $\hat{\Re}$  can be written as

$$\widehat{\mathfrak{P}} = (Q_{u_{0,1}}\mathcal{T} + Q_{u_{0,0}}) \circ \widehat{\mathfrak{R}} \circ (Q_{u_{0,1}}\mathcal{T} + Q_{u_{0,0}})^{-1}$$

**Corollary 1.** 1. Under the conditions of Theorem 1, the pseudo-difference operator  $\Re$  satisfies the following equations

$$\mathcal{T}(\mathfrak{R}^n) - \mathfrak{R}^n = [\Phi \circ \mathfrak{R}^n, \Phi^{-1}], \qquad n \in \mathbb{Z}, \qquad (2)$$

2. Under the conditions of Theorem 2, the pseudodifference operator  $\hat{\Re}$  satisfies equations

$$\mathcal{S}(\hat{\mathfrak{R}}^n) - \hat{\mathfrak{R}}^n = [\Psi \circ \hat{\mathfrak{R}}^n, \Psi^{-1}] \qquad n \in \mathbb{Z},$$

**Proof**. It that  $D_Q \circ \Re^n = \Re^n \circ D_Q, n \in \mathbb{Z}$ . Thus, we can apply Theorem 1 to  $\Re^n$  and  $\Re^n$  to produce (2). The proof of the second part of the Corollary is similar.

# Formal difference series, difference Adler Theorem.

**Definition 6.** A formal Laurent series of order N is defined as a formal semi-infinite sum

 $A = a_N \mathcal{S}^N + a_{N-1} \mathcal{S}^{N-1} + \dots + a_1 \mathcal{S} + a_0 + a_{-1} \mathcal{S}^{-1} + \dots,$ where  $a_k \in \mathcal{F}_Q$ ,  $a_N \neq 0$ ,  $N \in \mathbb{Z}$ .

A formal Taylor series of order -N is defined as a formal semi-infinite sum

 $C = c_{-N} \mathcal{S}^{-N} + c_{1-N} \mathcal{S}^{1-N} + \dots + c_{-1} \mathcal{S}^{-1} + c_0 + c_1 \mathcal{S} + \dots,$ where  $c_k \in \mathcal{F}_Q, \quad c_{-N} \neq 0, \quad N \in \mathbb{Z}.$ 

Laurent formal series form a skew-field. Sums and products (compositions) of formal series are formal series. The product is associative, but not commutative.

For any formal series A there exists a formal series  $A^{-1}$ such that  $A \circ A^{-1} = A^{-1} \circ A = 1$ . In order to find the first n coefficients of  $A^{-1}$  one needs to know exactly the first n coefficients of A.

Any pseudo-difference operator B can be uniquely represented by a Laurent formal series  $B_L$ . For example for  $B = (aS + b)^{-1}$  we have

$$B_L = \alpha_{-1} \mathcal{S}^{-1} + \alpha_{-2} \mathcal{S}^{-2} + \alpha_{-3} \mathcal{S}^{-3} + \cdots,$$

where the coefficients  $\alpha_k \in \mathcal{F}_Q$  can be found recursively:

$$\alpha_{-1} = \mathcal{S}^{-1}\left(\frac{1}{a}\right), \ \alpha_{-n} = -\mathcal{S}^{-1}\left(\frac{\alpha_{1-n}b}{a}\right).$$

**Definition 7.** Let  $A_L$  denote the Laurent series representations of a pseudo-differential operator A. Then ord  $A_L$  is called the **Laurent order** of A.

**Definition 8.** Let A be a formal series of order N  $A = a_N S^N + a_{N-1} S^{N-1} + \dots + a_1 S + a_0 + a_{-1} S^{-1} + \dots,$ The residue res(A) and logarithmic residue resln(A) are defined as

 $\operatorname{res}(A) = a_0, \quad \operatorname{res} \ln(A) = \ln(a_N).$ 

**Theorem 3.** Let  $A = a_N S^N + a_{N-1} S^{N-1} \cdots$  and  $B = b_M S^M + b_{M-1} S^{M-1} \cdots$  be two Laurent formal series of order N and M respectively. Then

 $\operatorname{res}[A, B] = (S - 1)(\sigma(A, B)),$ 

where  $\sigma(A, B) \in \mathcal{F}_Q$ 

$$\sigma(A,B) = \sum_{n=1}^{N} \sum_{k=1}^{n} \mathcal{S}^{-k}(a_{-n}) \mathcal{S}^{n-k}(b_n) - \sum_{n=1}^{M} \sum_{k=1}^{n} \mathcal{S}^{-k}(b_{-n}) \mathcal{S}^{n-k}(a_n).$$

## Canonical conservation laws: integrability conditions

**Theorem 4.** If a difference equation possesses a recursion operator  $\mathfrak{R}$ ,  $\operatorname{ord}_L(\mathfrak{R}) = N > 0$ , then it has infinitely many canonical conservation laws

$$(T-1)\rho_{nN} = (S-1)\sigma_{nN}, \quad n = 0, 1, 2, \dots$$

with canonical conserved densities

 $\rho_0 = \operatorname{res} \ln \mathfrak{R}_L, \quad \rho_{nN} = \operatorname{res} \mathfrak{R}_L^n, \quad n > 0,$ 

and the corresponding canonical fluxes

$$\sigma_0 = \sum_{k=0}^{N-1} \mathcal{S}^{k-1}(\ln \alpha_0), \quad \sigma_{nN} = \sigma(\Phi_L \circ \mathfrak{R}_L^n, \Phi_L^{-1})$$

where  $\alpha_0$  is the first coefficient in the Laurent expansion of  $\Phi$ :  $\alpha_0 = S^{-1} \left( \frac{Q_{u_{1,0}}}{Q_{u_{1,1}}} \right)$  and  $\sigma(A, B)$  is defined in Theorem 3.

If a recursion operator is known, then Theorem 4 gives us a completely algorithmic way to find explicitly a sequence of conservation laws, including both the densities  $\rho_k$  and the corresponding fluxes  $\sigma_k$ . The residues of powers a formal series are easy to compute. For instance, if

$$\Re = r_1 \mathcal{S} + r_0 + r_{-1} \mathcal{S}^{-1} + r_{-2} \mathcal{S}^{-2} + r_{-3} \mathcal{S}^{-3} + r_{-3} \mathcal{$$

then

$$\begin{aligned} \operatorname{res} \ln \mathfrak{R} &= \ln r_{1}, \quad \operatorname{res} \mathfrak{R} = r_{0}, \\ \operatorname{res} \mathfrak{R}^{2} &= \mathcal{S}^{-1}(r_{1})r_{-1} + r_{0}^{2} + r_{1}\mathcal{S}^{-1}(r_{-1}), \\ \operatorname{res} \mathfrak{R}^{3} &= \mathcal{S}^{-2}(r_{1})\mathcal{S}^{-1}(r_{1})r_{-2} + \\ &+ \mathcal{S}^{-1}(r_{0})\mathcal{S}^{-1}(r_{1})r_{-1} + 2\mathcal{S}^{-1}(r_{1})r_{-1}r_{0} + \mathcal{S}^{-1}(r_{1})r_{1}\mathcal{S}(r_{-2}) + \\ &+ r_{0}^{3} + 2r_{0}r_{1}\mathcal{S}(r_{-2}) + r_{1}\mathcal{S}(r_{-1})\mathcal{S}(r_{0}) + r_{1}\mathcal{S}(r_{1})\mathcal{S}^{2}(r_{-2}). \end{aligned}$$

**Proposition 1.** If a recursion operator  $\Re$  is represented by a first order formal series  $\Re_L = r_1 S + r_0 + r_{-1} S^{-1} + \cdots$ , then

(i) 
$$(\mathcal{T}-1)(\ln r_1) = (\mathcal{S}-1)\mathcal{S}^{-1}\left(\ln \frac{Q_{u_{1,1}}}{Q_{u_{1,0}}}\right),$$

(ii) 
$$(\mathcal{T}-1)(r_0) = (\mathcal{S}-1)\mathcal{S}^{-1}(r_1F),$$
  
(iii)  $(\mathcal{T}-1)(r_{-1}\mathcal{S}^{-1}(r_1) + r_0^2 + r_1\mathcal{S}(r_{-1})) = (\mathcal{S}-1)(\sigma_2),$ 

where

$$\sigma_2 = S^{-1}(r_1 F) \left\{ S^{-1}(r_0) + r_0 - S^{-2}(r_1 F) \right\} - (1 + S^{-1}) \left( r_1 G S^{-1}(r_1 F) \right),$$

and F, G denote

$$F = \frac{Q_{u_{0,1}} \mathcal{S}^{-1}(Q_{u_{1,0}}) - Q_{u_{0,0}} \mathcal{S}^{-1}(Q_{u_{1,1}})}{Q_{u_{1,0}} \mathcal{S}^{-1}(Q_{u_{1,1}})}, \qquad G = \frac{Q_{u_{0,0}}}{Q_{u_{1,0}}}.$$

**Proposition 2.** If a recursion operator  $\hat{\Re}$  is represented by a first order formal Laurent  $\mathcal{T}$  series  $\hat{\Re}_L = \hat{r}_1 \mathcal{T} + \hat{r}_0 + \hat{r}_{-1} \mathcal{T}^{-1} + \cdots$ , then

(i) 
$$(S-1)(\ln \hat{r}_1) = (T-1)T^{-1}\left(\ln \frac{Q_{u_{1,1}}}{Q_{u_{0,1}}}\right),$$

(ii) 
$$(S-1)(\hat{r}_0) = (T-1)T^{-1}(\hat{r}_1\hat{F}),$$

(iii) 
$$(\mathcal{S}-1)(\hat{r}_{-1}\mathcal{S}^{-1}(\hat{r}_1) + \hat{r}_0^2 + \hat{r}_1\mathcal{S}(\hat{r}_{-1})) = (\mathcal{T}-1)(\hat{\sigma}_2),$$

where

$$\hat{\sigma}_{2} = S^{-1}(\hat{r}_{1}\,\hat{F}) \left\{ S^{-1}(\hat{r}_{0}) + \hat{r}_{0} - S^{-2}\left(\hat{r}_{1}\,\hat{F}\right) \right\} - (1 + S^{-1}) \left(\hat{r}_{1}\,\hat{G}\,S^{-1}\left(\hat{r}_{1}\,\hat{F}\right)\right),$$

and F, G denote

$$\widehat{F} = \frac{Q_{u_{1,0}}\mathcal{T}^{-1}(Q_{u_{0,1}}) - Q_{u_{0,0}}\mathcal{T}^{-1}(Q_{u_{1,1}})}{Q_{u_{0,1}}\mathcal{T}^{-1}(Q_{u_{1,1}})}, \qquad \widehat{G} = \frac{Q_{u_{0,0}}}{Q_{u_{0,1}}}.$$

## The Viallet equation (Q5)

$$Q := a_1 u_{0,0} u_{1,0} u_{0,1} u_{1,1} + a_2 (u_{0,0} u_{1,0} u_{0,1} + u_{1,0} u_{0,1} u_{1,1} \\ + u_{0,1} u_{1,1} u_{0,0} + u_{1,1} u_{0,0} u_{1,0}) \\ + a_3 (u_{0,0} u_{1,0} + u_{0,1} u_{1,1}) + a_4 (u_{1,0} u_{0,1} + u_{0,0} u_{1,1}) \\ + a_5 (u_{0,0} u_{0,1} + u_{1,0} u_{1,1}) \\ + a_6 (u_{0,0} + u_{1,0} + u_{0,1} + u_{1,1}) + a_7 \\ = 0,$$

where  $a_i$  are free complex parameters.

- All of the ABS (Adler, Bobenko, Suris) equations can be obtained from  ${\rm Q}_5;$
- For generic coefficients it can be transformed into Adler's equation  $(\mathbf{Q}_4)$ ;
- $\mathbf{Q}_5$  is invariant under  $u_{1,0} \leftrightarrow u_{0,1}$  and  $a_3 \leftrightarrow a_5$ .

Genralised symmetries K:  $D_Q[K] = 0$ 

- Tongas-Tsoubelis-Xenitidis; Rasin-Hydon: the ABS equations (mastersymmetries);
- A generalised symmetry of order (-1,1) for the Viallet equation (Tongas-Tsoubelis-Xenitidis):

$$K^{(1)} := \frac{h}{u_{1,0} - u_{-1,0}} - \frac{1}{2} \partial_{u_{1,0}} h = \frac{h_{-1}}{u_{1,0} - u_{-1,0}} + \frac{1}{2} \partial_{u_{-1,0}} h_{-1},$$
  
where  $h(u_{0,0}, u_{1,0}) = Q \partial_{u_{0,1}} \partial_{u_{1,1}} Q - \partial_{u_{0,1}} Q \partial_{u_{1,1}} Q.$ 

Statement: The Viallet equation possesses a symmetry of order (-2,2)

$$K^{(2)} = \frac{h h_{-1}}{(u_{1,0} - u_{-1,0})^2} \left( \frac{1}{u_{2,0} - u_{0,0}} + \frac{1}{u_{0,0} - u_{-2,0}} \right).$$

### **Recursion operator of the Viallet equation**

**Theorem**. The Viallet equation possesses a recursion operator  $\mathcal{R} = \mathcal{H} \circ \mathcal{I}$ .

$$\begin{split} \mathcal{I} &= \frac{1}{h} \mathcal{S} - \mathcal{S}^{-1} \frac{1}{h} \text{ ; (symplectic operator)} \\ \mathcal{H} &= \frac{h_{-1} h h_1}{(u_{1,0} - u_{-1,0})^2 (u_{2,0} - u_{0,0})^2} \mathcal{S} \\ &- \mathcal{S}^{-1} \frac{h_{-1} h h_1}{(u_{1,0} - u_{-1,0})^2 (u_{2,0} - u_{0,0})^2} \\ &+ 2 \, K^{(1)} \, \mathcal{S} \, (\mathcal{S} - 1)^{-1} K^{(2)} + 2 \, K^{(2)} \, (\mathcal{S} - 1)^{-1} K^{(1)} \\ & \text{(Hamiltonian operator)} \ . \end{split}$$

## Symplectic and Hamiltonian Operators

$$\mathcal{I}:\mathfrak{h}\to\Omega^1\Longrightarrow\omega=rac{1}{2}\int\!\mathrm{d} u\wedge\mathcal{I}\mathrm{d} u$$

if 2-form  $\omega$  is anti-symmetric and is closed, we say  $\mathcal{I}$  is **symplectic**.

$$\mathcal{H}: \Omega^1 \to \mathfrak{h} \Longrightarrow \left\{ \int f, \int g \right\} = <\delta(f), \mathcal{H}\delta(g) >$$

if the Poisson bracket defined is anti-symmetric and satisfies the Jacobi identity, we say  $\mathcal{H}$  is **Hamiltonian**.

(cf. Dorfman, 1993 & Olver 1993)

### Symplectic and Hamiltonian Operators of Q5

Symplectic: symmetries  $\mapsto$  covariants Hamiltonian: covariants  $\mapsto$  symmetries

$$\begin{aligned} \mathcal{I}(K^{(1)}) &= \frac{1}{u_{2,0} - u_{0,0}} - \frac{1}{u_{0,0} - u_{-2,0}} + \frac{1}{2} \frac{\partial_{u_{0,0}} h}{h} + \frac{1}{2} \frac{\partial_{u_{0,0}} h_{-1}}{h_{-1}} \\ &= \delta_{s} \left( \ln h_{-1} - 2 \ln(u_{1,0} - u_{-1,0}) \right) \\ \mathcal{HI}(K^{(1)}) &= \frac{h h_{-1}}{(u_{1,0} - u_{-1,0})^{2}} \left( \frac{K_{2}^{(1)}}{(u_{2,0} - u_{0,0})^{2}} + \frac{K_{-2}^{(1)}}{(u_{0,0} - u_{-2,0})^{2}} \right) \\ &+ \left( \frac{1}{u_{2,0} - u_{0,0}} + \frac{1}{u_{0,0} - u_{-2,0}} \right) K^{(1)} K^{(2)} \end{aligned}$$

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